

LIGO

**COSMIC
EXPLORER**



Terrestrial GW Detectors, present and future

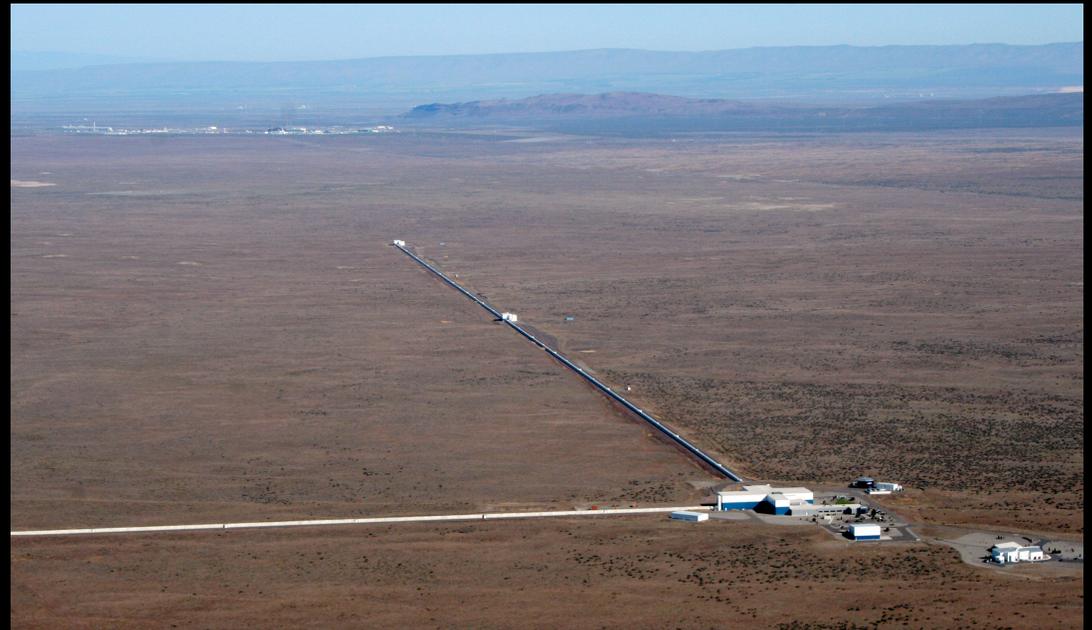
David Shoemaker

MIT Kavli Institute – LIGO Laboratory

Slides from many CE and LIGO Collaborators!

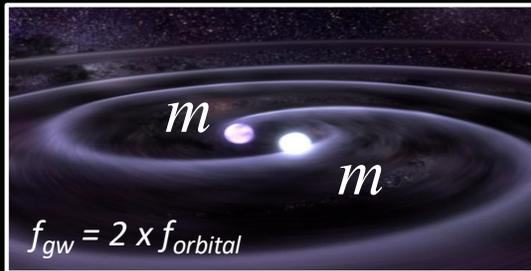
GWJ 2025 Colloquium

3 June 2025
CE-G2500043



Properties of GR's Gravitational Waves

Binary Coalescence of two compact objects



GW generation:
lowest order radiation is quadrupole

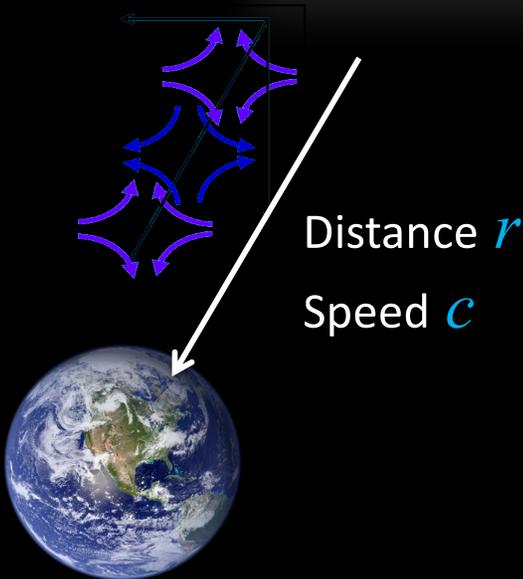
metric perturbation \rightarrow $h = \frac{2G}{c^4 r} \ddot{I}$ \leftarrow quadrupole moment

Two masses m in a circular orbit at a distance r create a periodic strain h in space

$$h = \frac{2Gm}{c^4 r} (2\pi f_{gw})^{2/3}$$

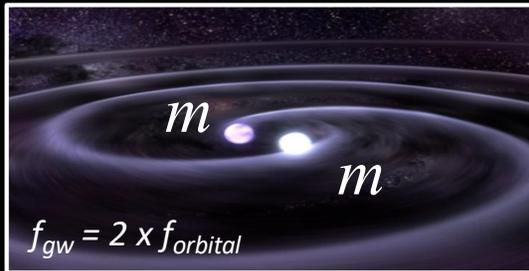
About once a week,
a wave passes with this characteristic strain:

$$1.5 \times 10^{-21} \left(\frac{m}{30M_{\odot}} \right) \left(\frac{400 \text{ Mpc}}{r} \right) \left(\frac{f_{gw}}{50 \text{ Hz}} \right)^{2/3}$$



Gravitational Wave Properties

Binary Coalescence of two compact objects



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lowest order radiation is quadrupole

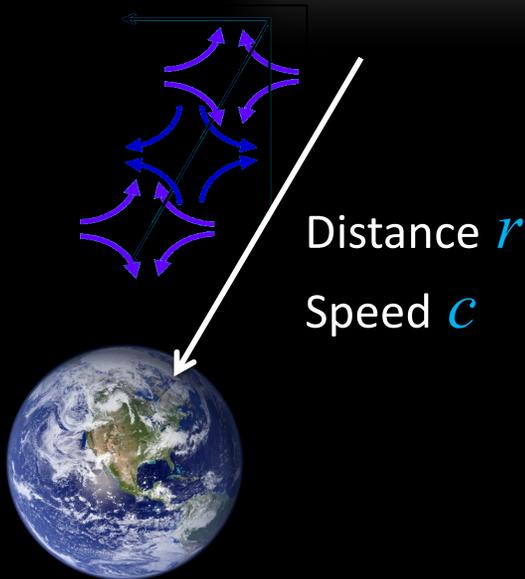
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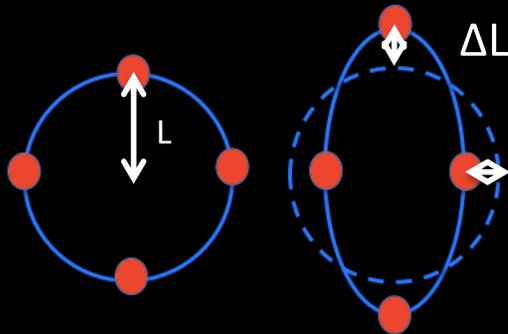
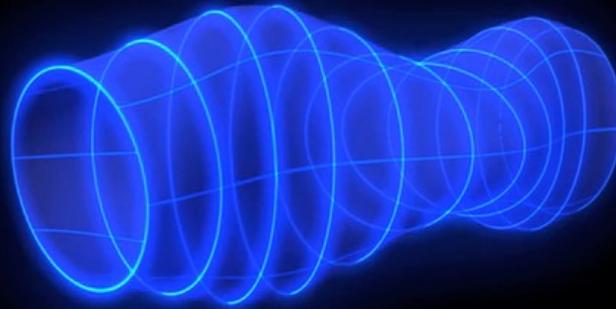
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Stretching and squeezing of space-time



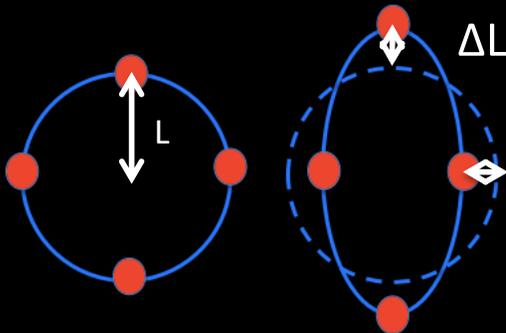
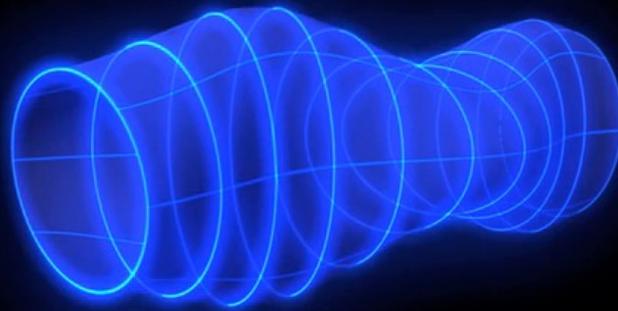
Amplitude of the gravitational wave strain is $h = \Delta L/L$

$$\Delta L = h L$$

Big L makes ΔL easier to measure; current detectors have $L = 4 \text{ km}$, so from our two-mass example

$$\sim 10^{-21} \times \sim 10^3 = \sim 10^{-18} \text{ m}$$

Stretching and squeezing of space-time



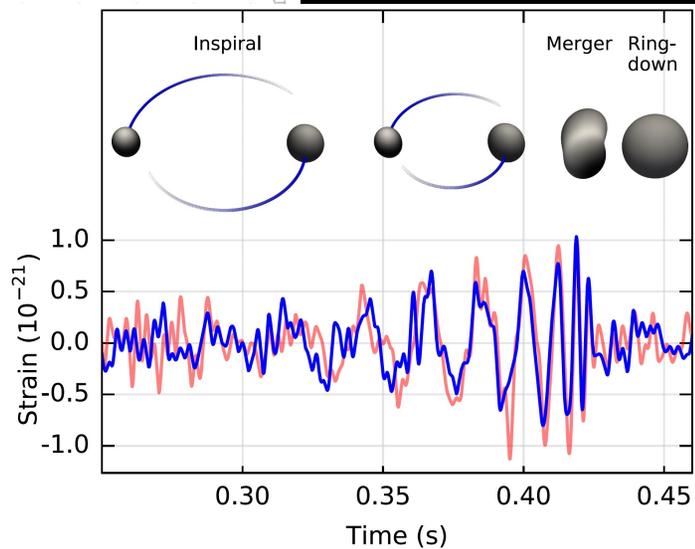
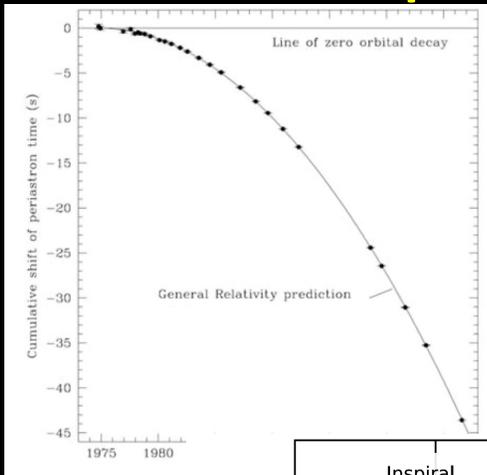
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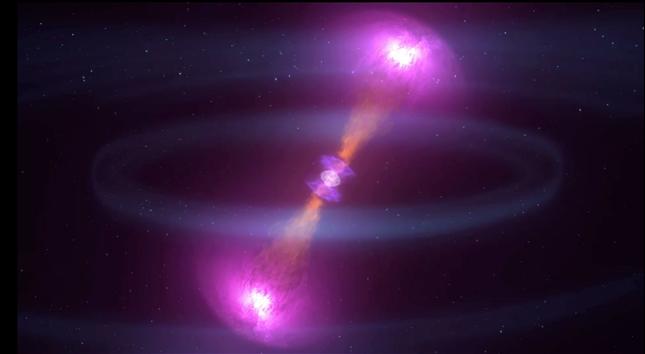
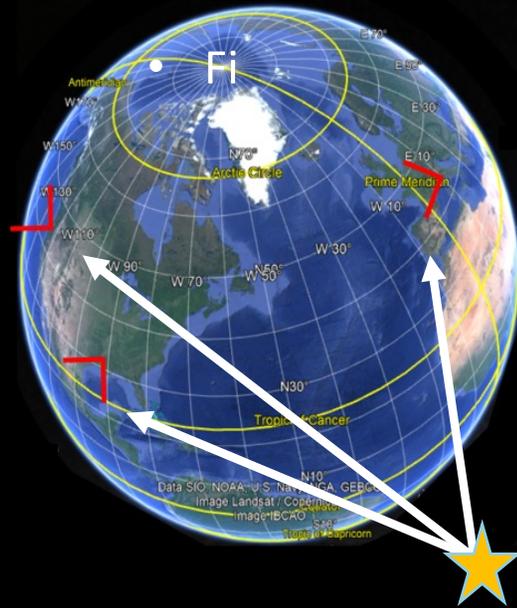
On September 14, 2015 at 09:50:45 UTC



GW150914

- First direct observation of the time-varying strain due to GWs
 - Hulse and Taylor deserved their Nobel!
- Two black holes, $36 M_{\odot}$ and $29 M_{\odot}$, merged merged 1.3 billion years ago to form a single black hole of $62 M_{\odot}$
- As they merged, the equivalent of $3 M_{\odot}$ was emitted in GW in ~ 0.2 sec
- Was seen in the two **LIGO** detectors separated by 3000 km, with a time difference consistent with the speed of light

The VIRGO detector, located in Italy, joined LIGO in August 2017 – and we observed a Binary Neutron Star coalescence a few weeks later

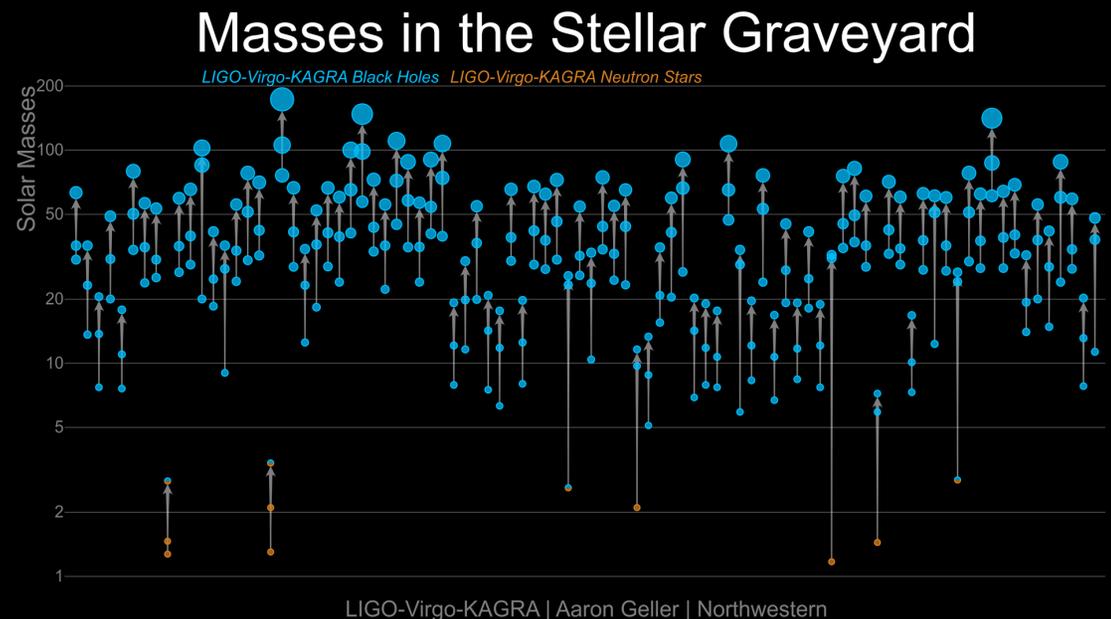


GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral. LVC (2017) PRL 119, 161101

- Made first joint GW-EM source observation
- Linked short GRBs, binary neutron stars, kilonovae
- Independently measured the local Hubble constant
- Measured the speed of gravitational-wave propagation
- Made initial constraints on the neutron star equation of state
- Constrained the rate of binary neutron star mergers in the local Universe (and thus their production of heavy metals)

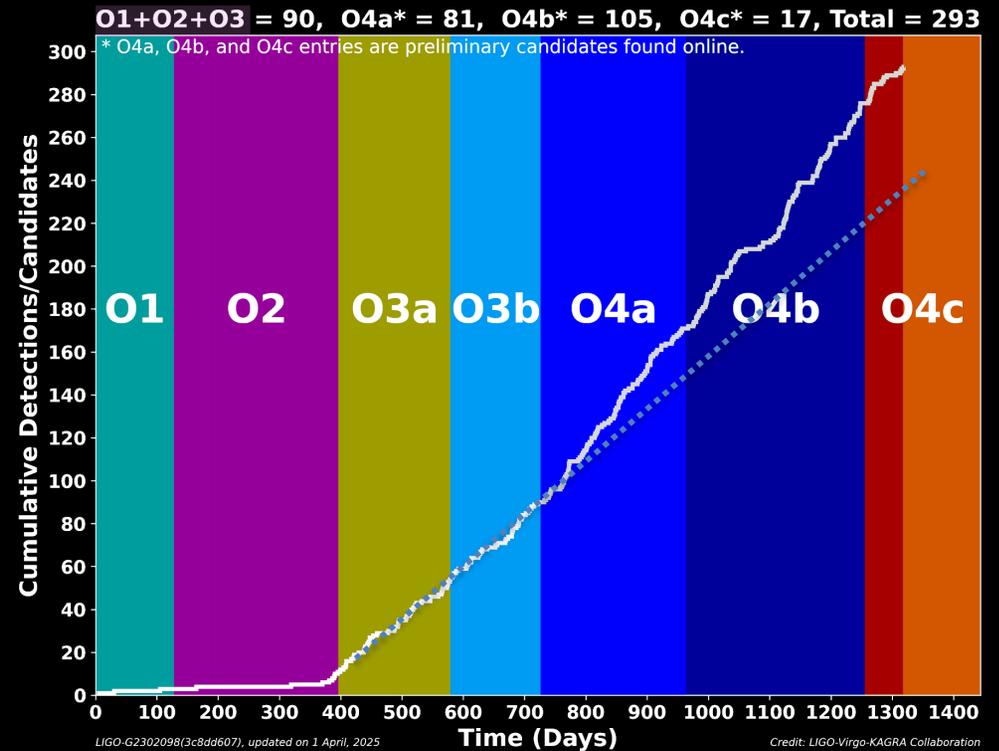
Astrophysics, Cosmology, Testing GR...

- Binaries: population studies of BH, NS
- Hubble tension resolution, other measures of large-scale space-time
- Nuclear physics – Neutron-star EOS
- Searches for deviations from GR
- Sub-solar BH of primordial origin
- Searches for Exotic matter
- Sources other than Binaries
 - Supernovae
 - Pulsars
 - Strings, Cusps, et alia
 - Stochastic backgrounds, both astrophysical and cosmological
- **Just getting started**

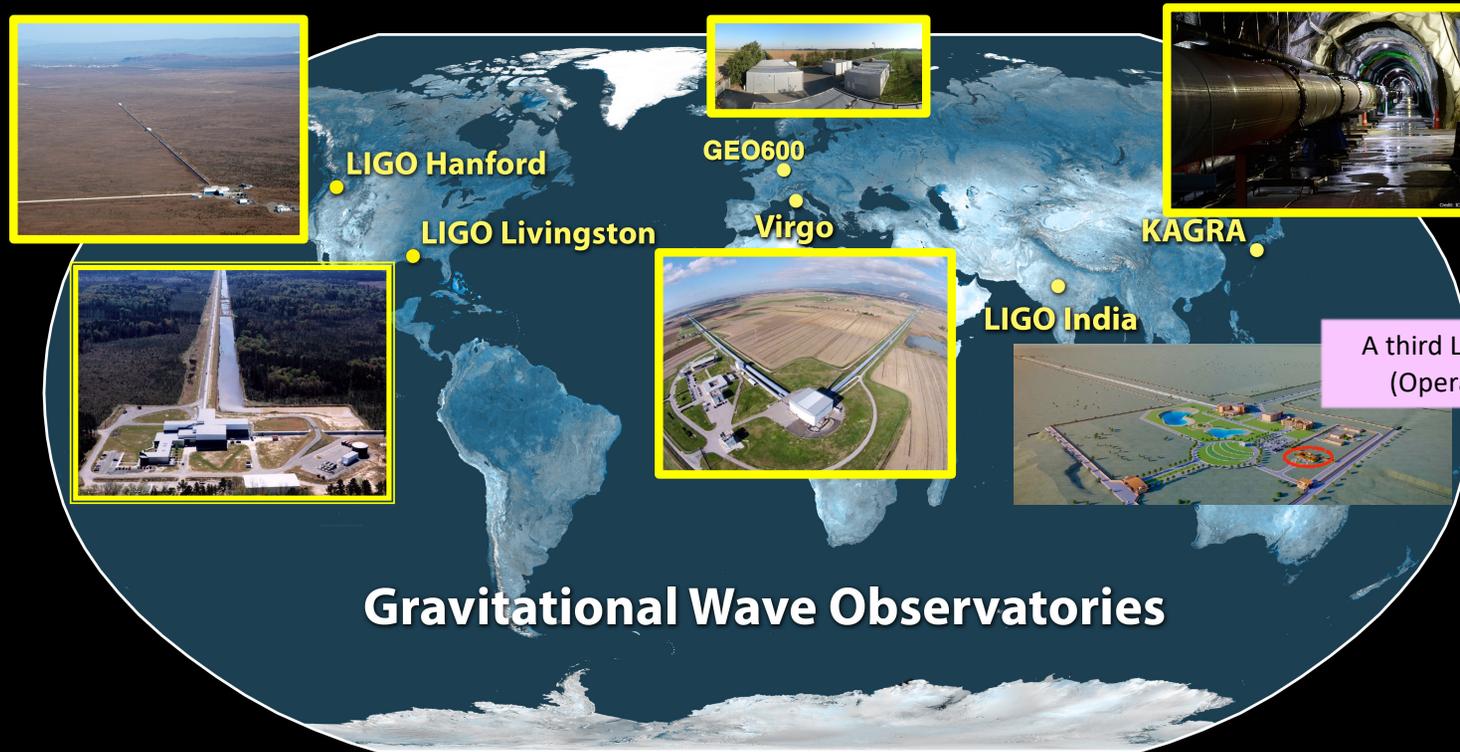


Sensitivity improvements boost signal rate

- Volume of space – and thus # of target sources – grows as $\sim(\text{sensitivity})^3$
- There are gaps for the upgrades... but so far we make up for it



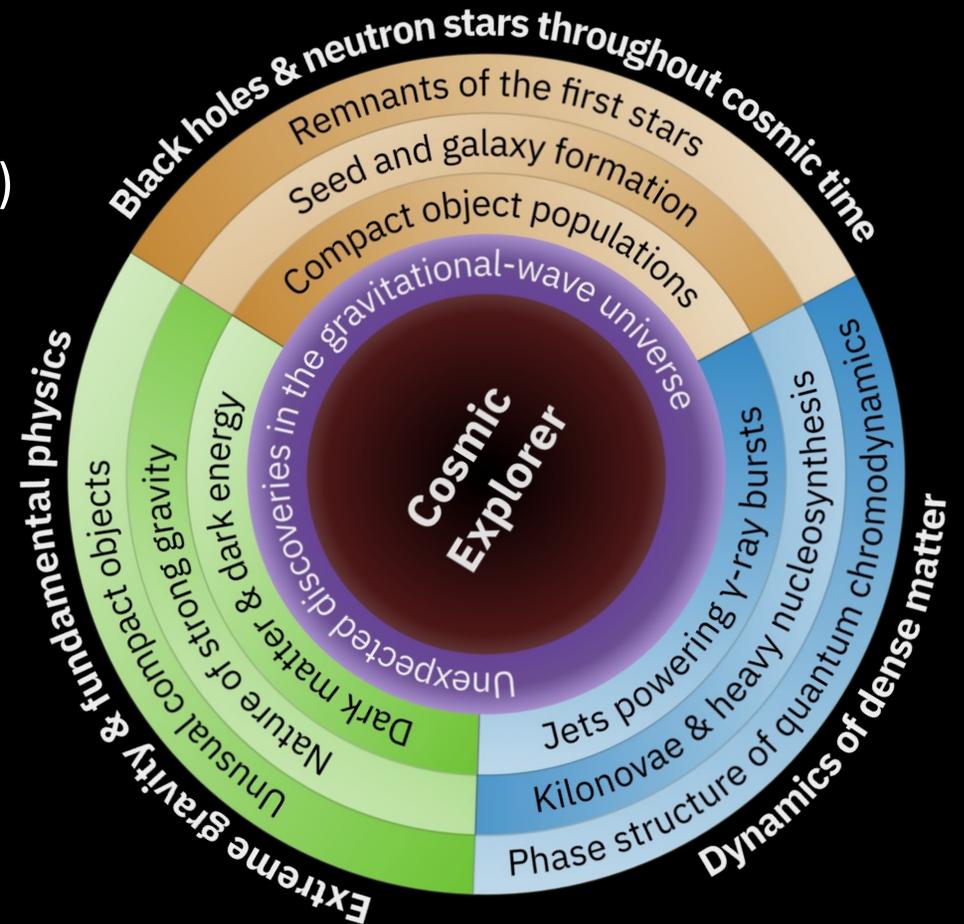
The current ground-based gravitational-wave world-wide network



Gravitational Wave Observatories

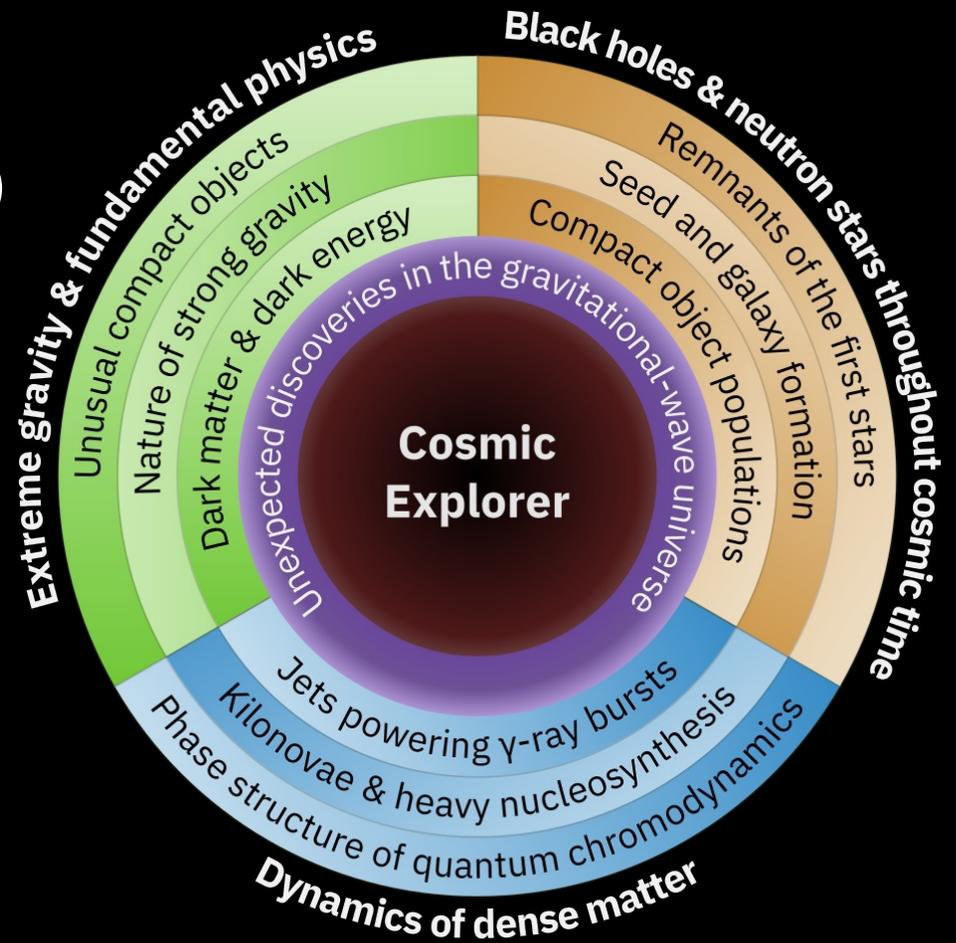
What could we do with 10x better GW detectors?

- Event rate boost
(10x sensitivity \rightarrow 10^5 sources per year)
- Increases resolution of waveforms
 - tests of GR, BH Ringdown
 - Precise spin etc. measurements
- Wider bandwidth to expose Neutron star coalescence
 - dynamics of dense matter
- More Multi-Messenger Astrophysics
 - Where are those neutron stars....



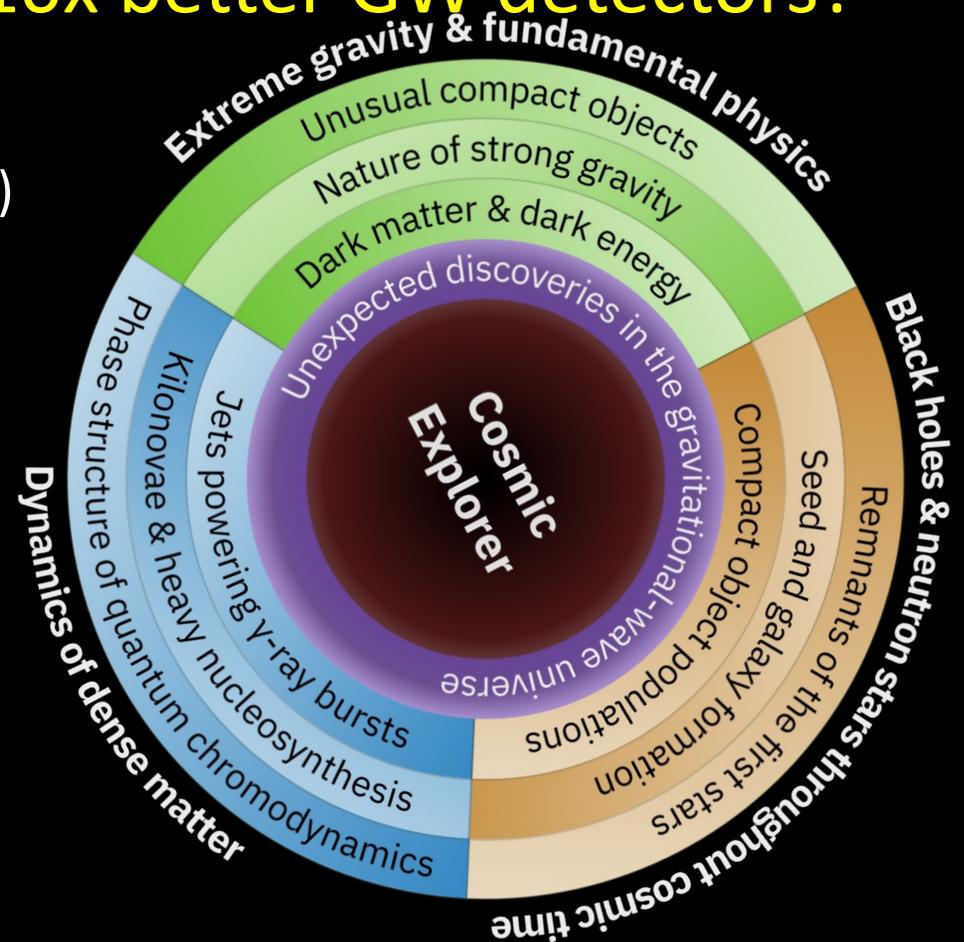
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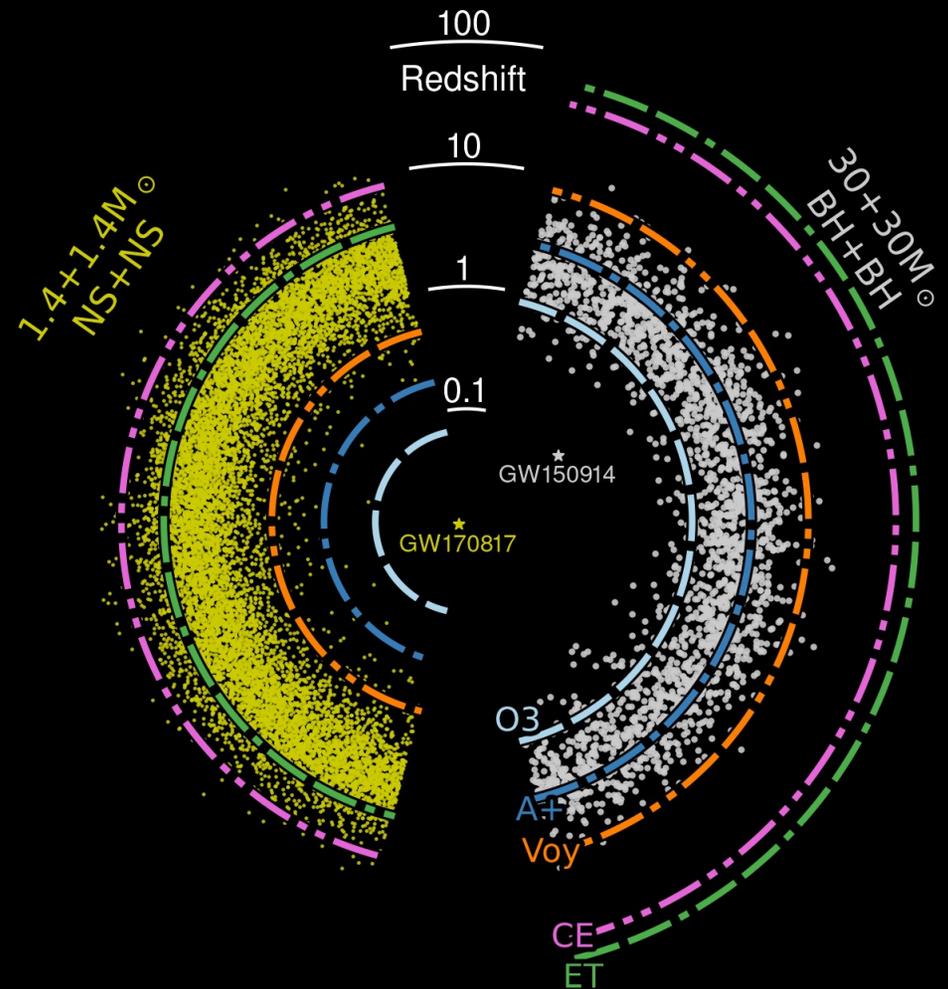
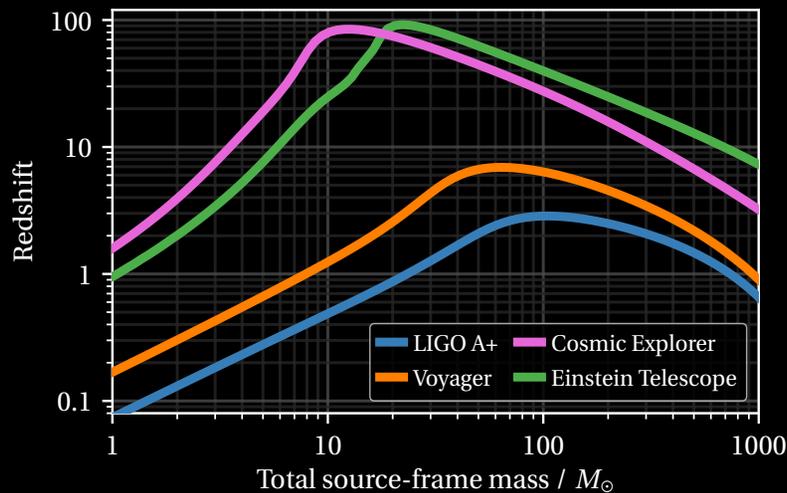
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Black Holes and Neutron Stars throughout cosmic time

- The best understood source of gravitational wave emissions are compact binary systems.
- Can build a detector able see all binaries in a broad range of masses

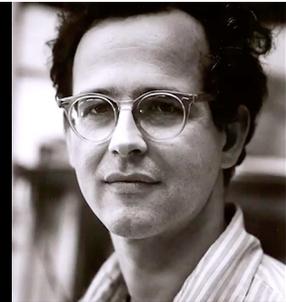


Even better detectors would deliver more science.
How to build a such a 10x better detector?

- **Current detector concept**
- Scaling laws

Basic principle for detecting gravitational waves: a laser Michelson interferometer

- A transducer from strain to light intensity \rightarrow electrical signal



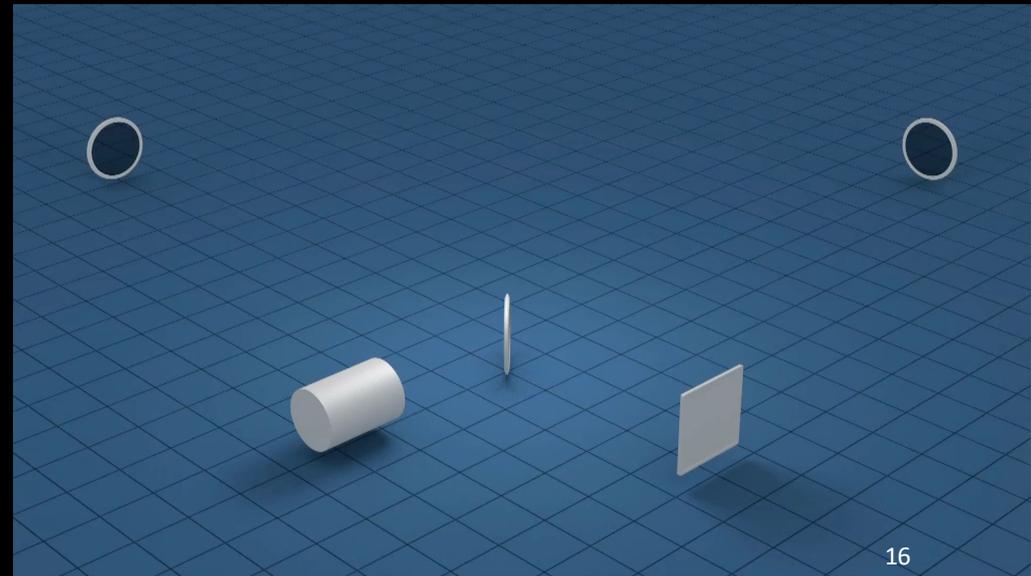
Rai Weiss,
early '70s

What determines the sensitivity of this detector?

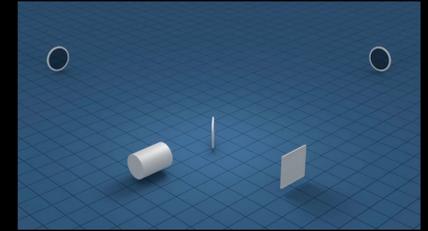
- **Distance over which strain is sensed**

$$\Delta L = h L$$

- Resolution of the optical sensing
- Stochastic forces making motion of mirrors, masking GWs

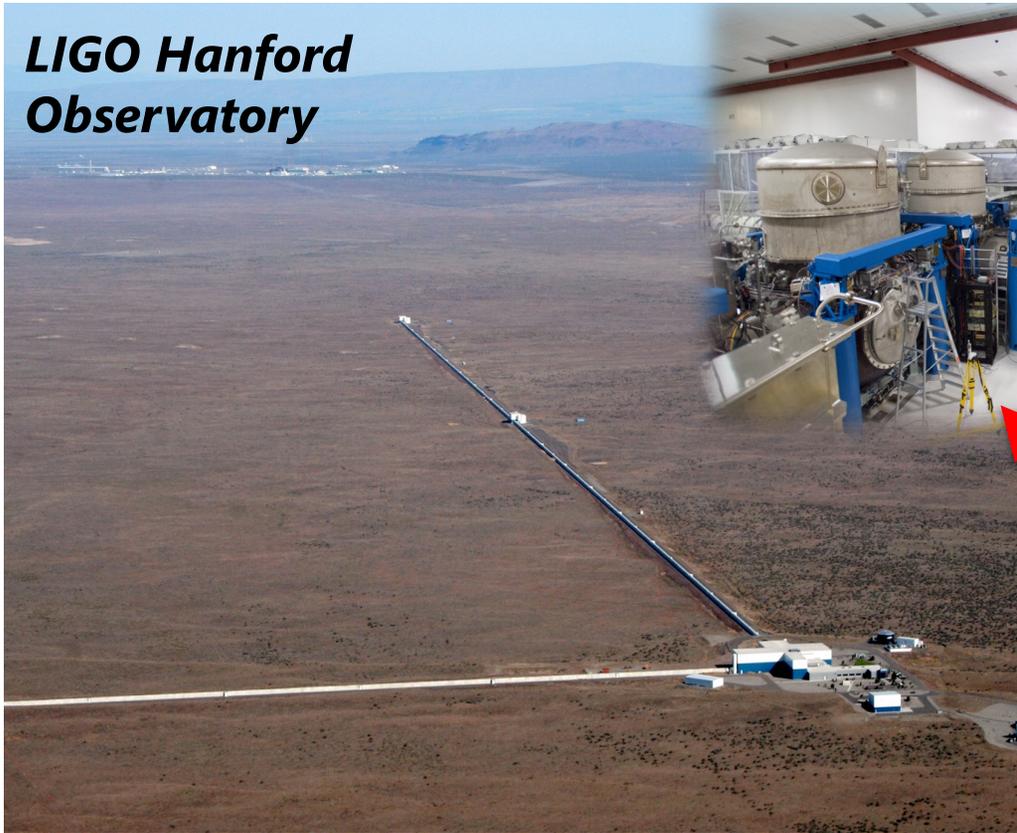


The Infrastructure for a realistic implementation



- Amplitude of the gravitational wave strain is $h = \Delta L/L \rightarrow \Delta L = h L$
- A sufficient length L of the arms is needed to bring the GW-induced strain to a measurable level (LIGO: 4km)
- Sensing laser light must travel in an excellent vacuum (LIGO: 10^{-9} Torr)
- The vacuum system diameter must accommodate a diffraction limited beam over 4km (LIGO: 1.2 m Diameter; $\rightarrow 10,000 \text{ m}^3$)
- The vacuum system must be *straight*, level, and protected from the human and natural environment (LIGO: earthmoving, concrete bed, aligned to several mm over 4km, and protected by a concrete cover)
- Corner and end buildings with particulate, temperature control; staff buildings; outreach/public science building (LIGO: $10,000 \text{ m}^2$)

**LIGO Hanford
Observatory**



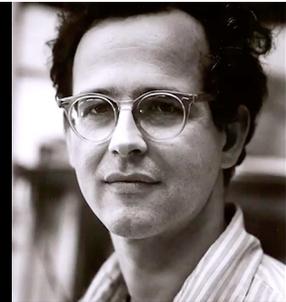
**Present LIGO Observatory sites
 $L = 4\text{km}$**

**LIGO Livingston
Observatory**



Basic principle for detecting gravitational waves: a laser Michelson interferometer

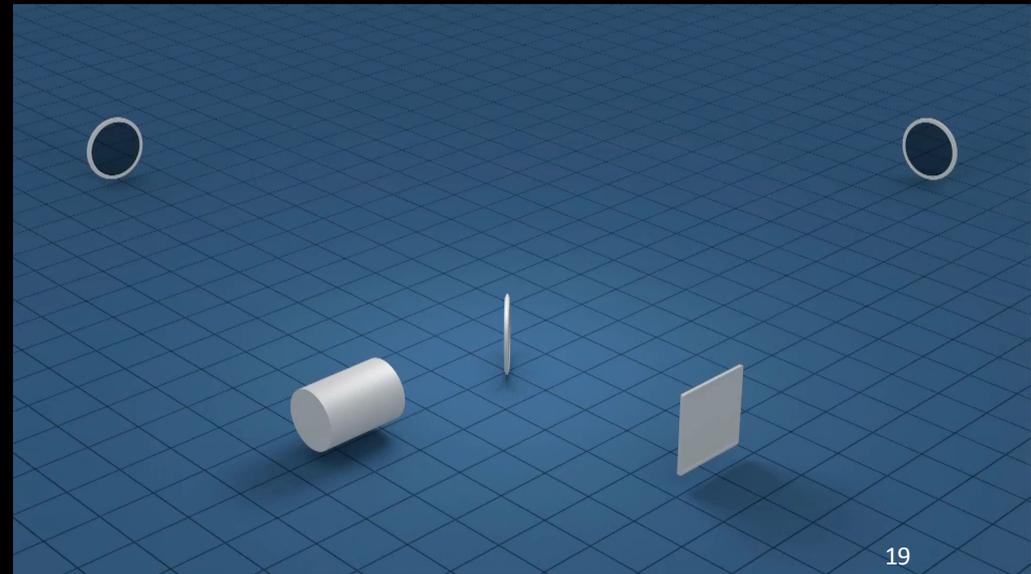
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- **Resolution of the optical sensing**
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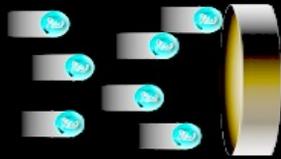


Resolution of the optical sensing

- **Shot noise** – ability to resolve a fringe shift due to a GW (counting statistics; *A. Einstein, 1909*)

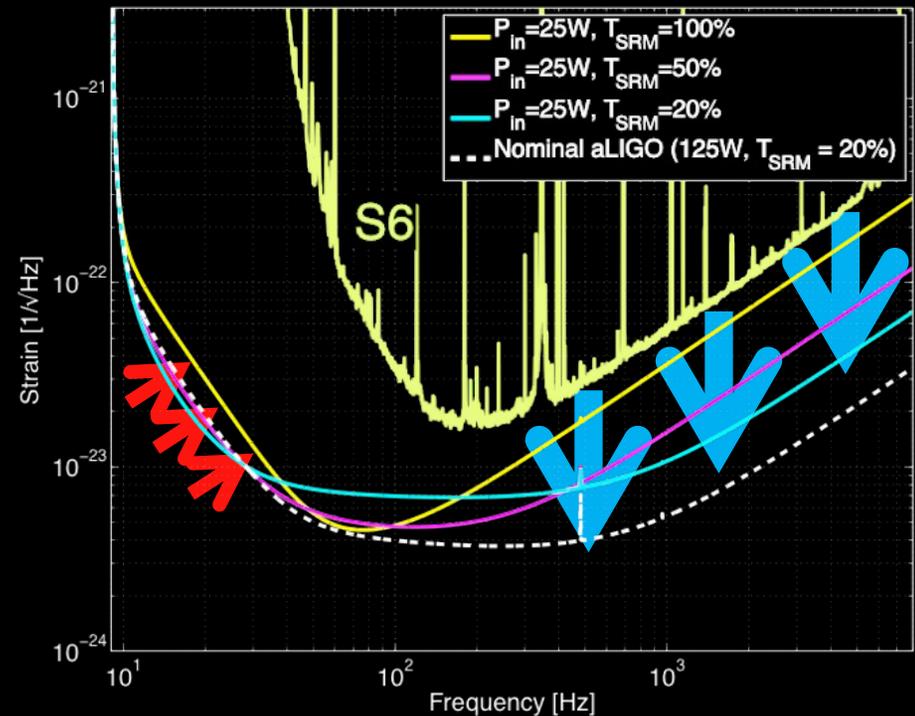
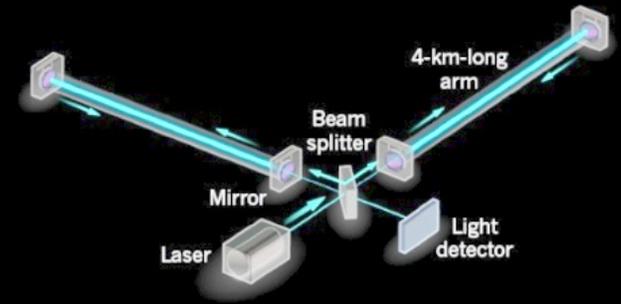
$$h_{\text{sn}}(f) = \frac{1}{L} \sqrt{\frac{\hbar c \lambda}{2\pi P}}$$

- **Radiation Pressure noise** – buffeting of test mass by photons increases low-frequency noise – use heavy test masses!

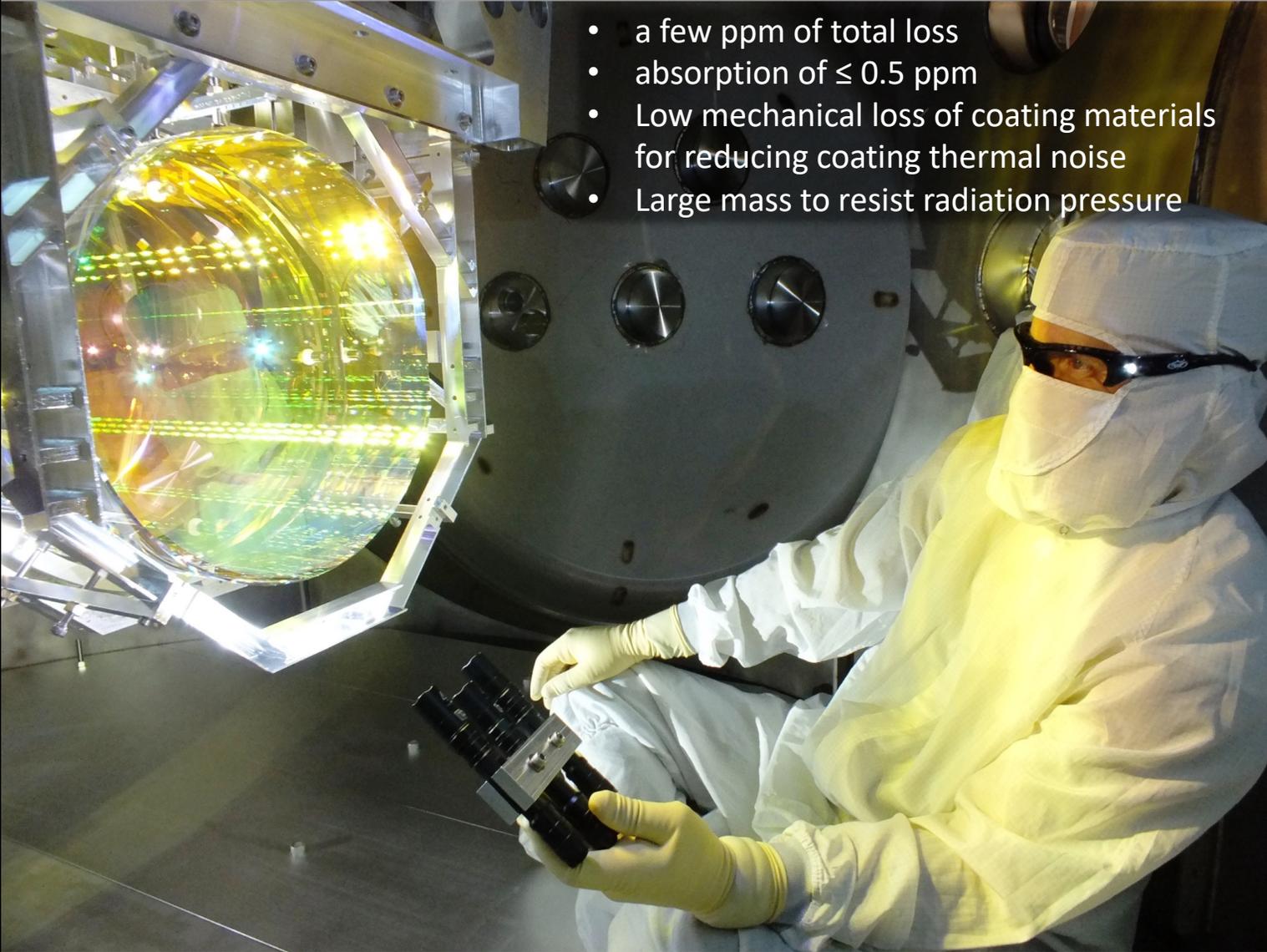


$$h_{\text{rp}}(f) = \frac{1}{m f^2 L} \sqrt{\frac{\hbar P}{2\pi^3 c \lambda}}$$

- **'Standard Quantum Limit'**



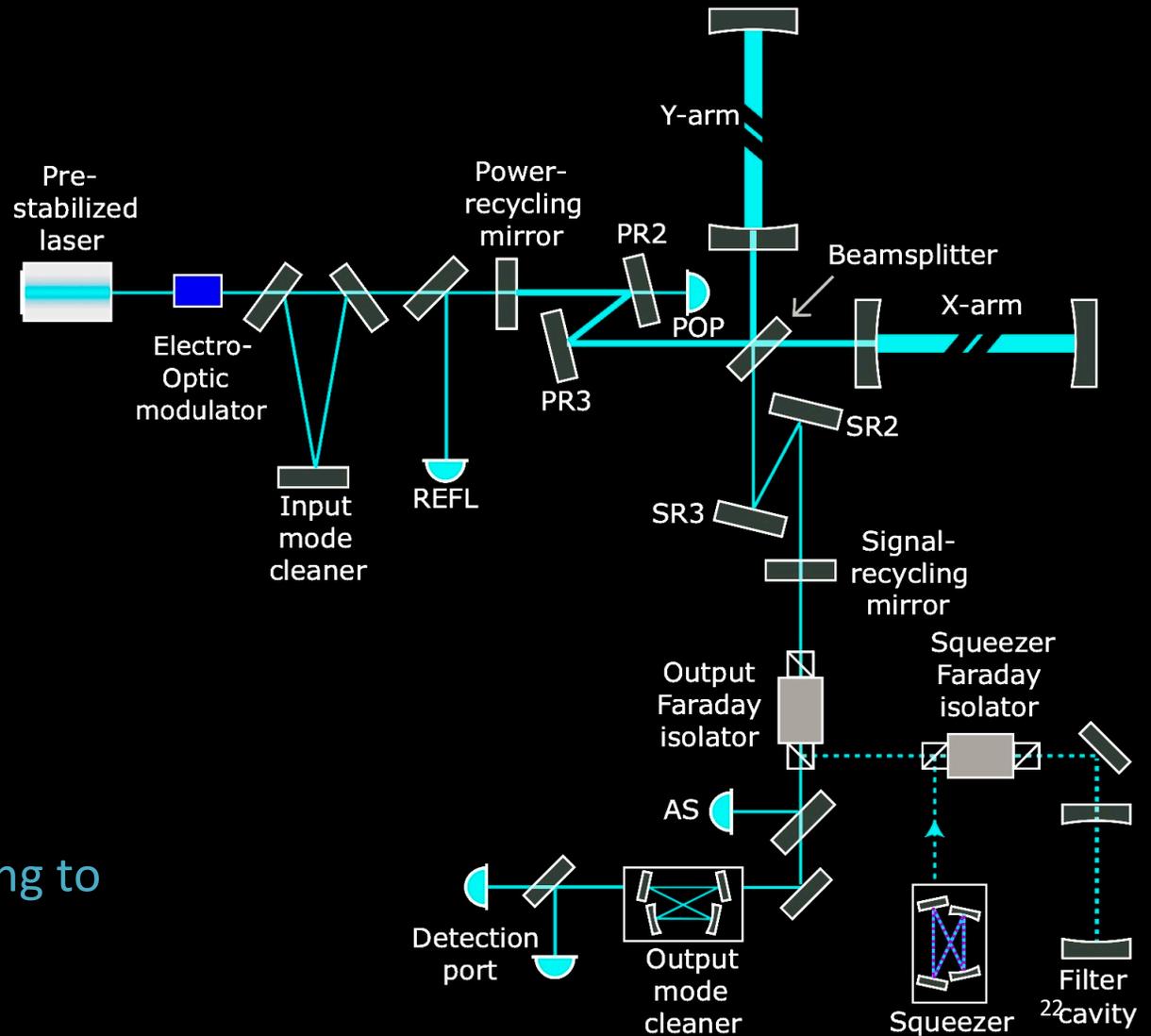
High Quality Optics



- a few ppm of total loss
- absorption of ≤ 0.5 ppm
- Low mechanical loss of coating materials for reducing coating thermal noise
- Large mass to resist radiation pressure

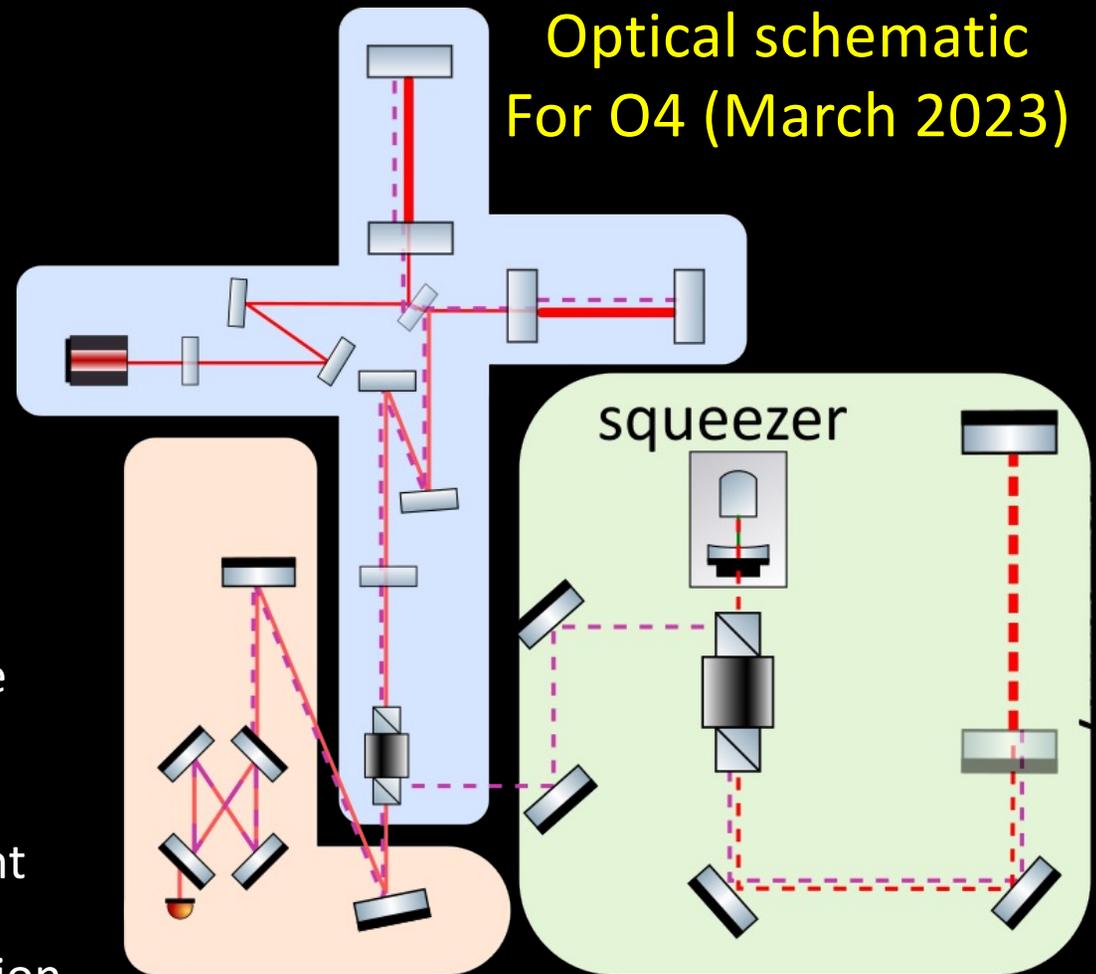
Optical schematic

- Fabry-Perot cavities in arms to increase interaction time
- Coupled cavities to increase power, extract GW frequencies
- FP cavities in transmission to enforce TEM₀₀ mode
- ~100 degrees-of-freedom to control!
- Frequency-dependent squeezing to manage quantum noise

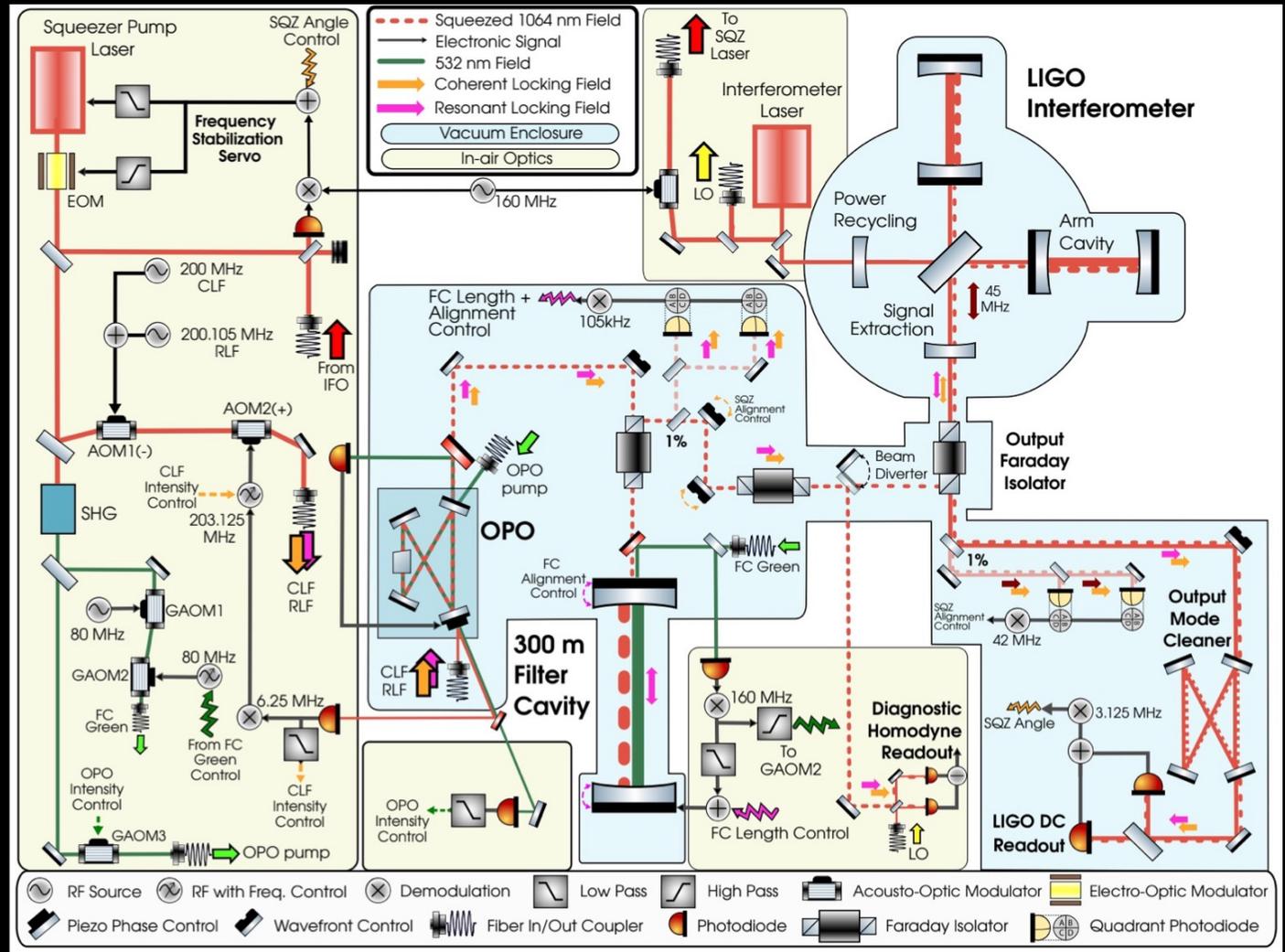


Squeezing

- Heisenberg Uncertainty Principle dictates that precise values of phase, and amplitude, of light cannot be known at the same time
- $\Delta X_{\text{phase}} \Delta Y_{\text{amp}} \leq h_{\text{Heisenberg}}/2$
- We can choose however to e.g., know the amplitude less well and look more closely at the phase
- ‘Squeezed light’ used in O3 to reduce shot noise at the expense of more radiation pressure noise
- For O4, adding ‘Frequency Dependent Squeezing’ to reduce high-frequency shot noise, and low-frequency radiation pressure noise



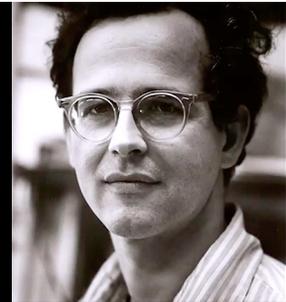
That Squeezer in a bit more detail...



Ganapathy*,Jia*,Nakano*,Xu*,LIGO O4 instrument scientists,
Broadband Quantum Enhancement of the LIGO Detectors with Frequency-Dependent Squeezing, PRX 13, 041021 (2023)

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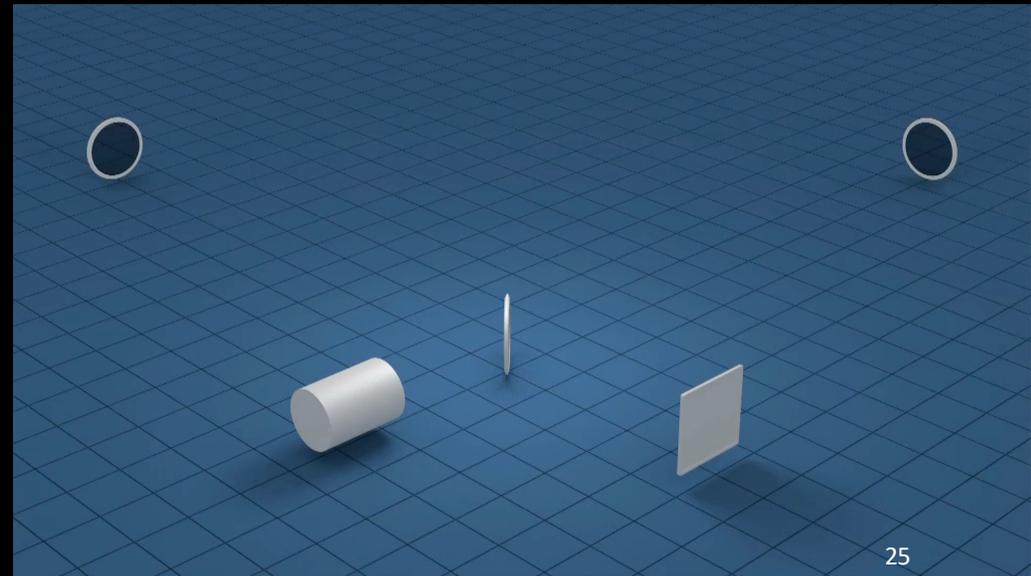
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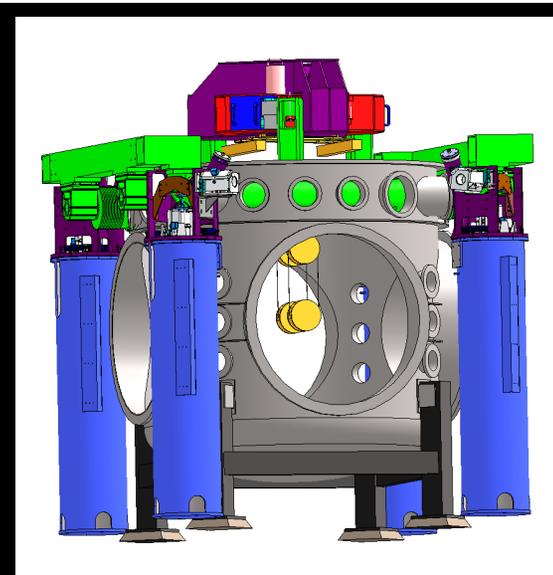
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- Resolution of the optical sensing
- **Stochastic forces making motion of mirrors, masking GWs**



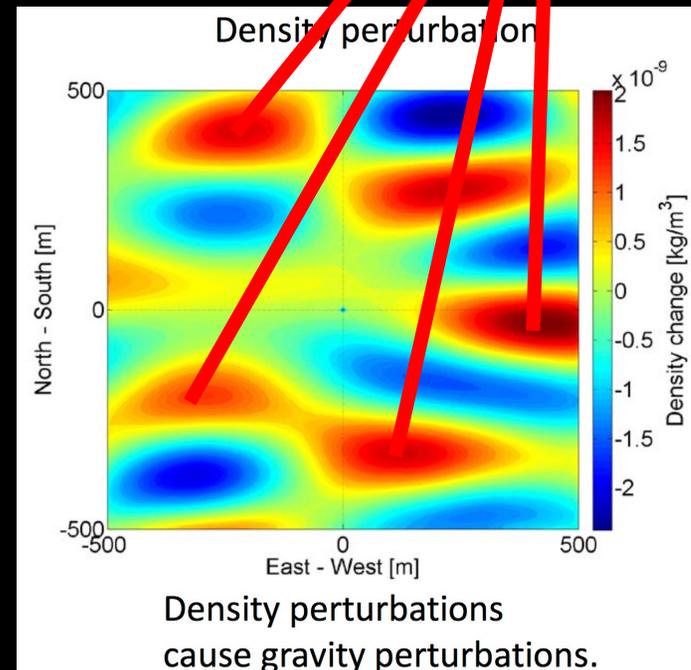
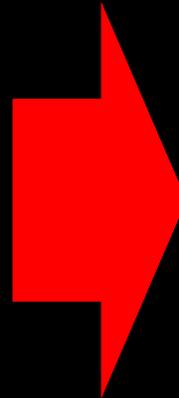
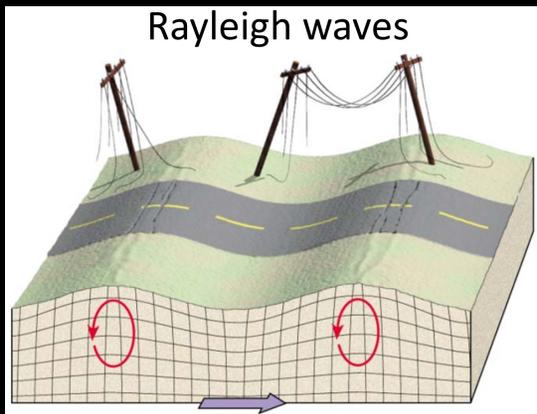
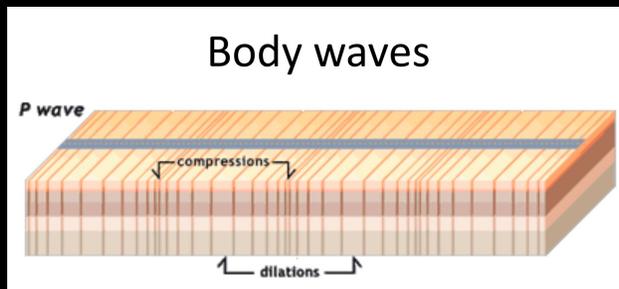
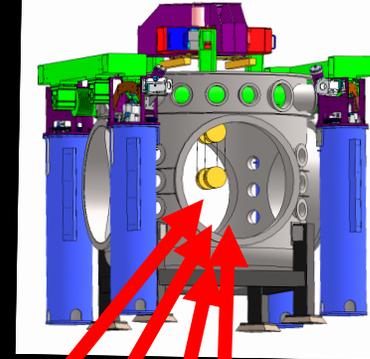
Stochastic force motion of mirrors

- Suppress physical transmission of (seismic etc.) motion from the outside via:
 - 3x6 degree-of-freedom servo controlled platforms
 - 4 pendulums in series, giving $((f_o/f)^2)^4$ isolation
 - Order of 10^{12} suppression of motion
- Newtonian gravitational gradients due to seismic activity
 - Limit on lowest frequency for ground-based detectors...
→ go to space (LISA)
- Concentrate thermal noise in frequency bands to limit impact on GW range



Newtonian background due to seismic environment

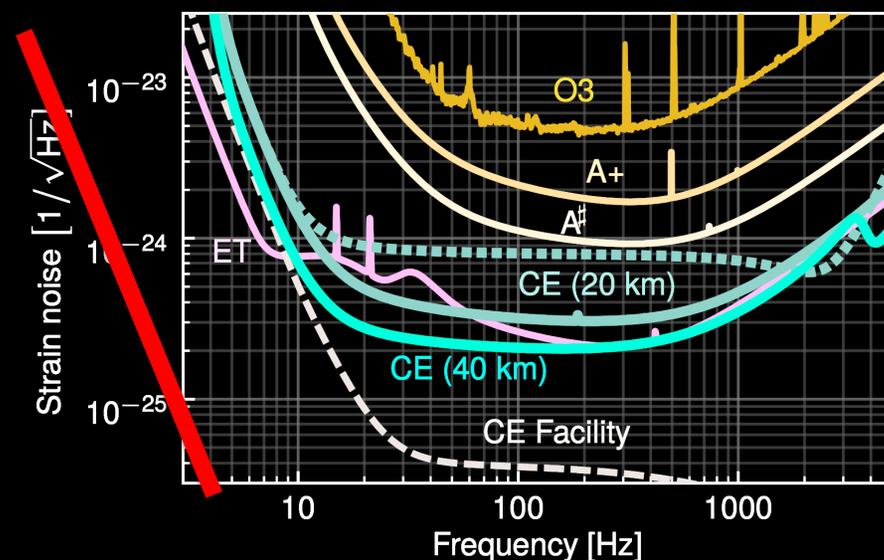
- In principle, can eliminate all direct mechanical coupling
- Can not shield against the wandering net gravity vector



Images: J. Harms

Newtonian background creates a wall at a few Hz

- Newtonian Background falls as $\sim 1/f^5$
- Can reduce somewhat by moving underground
 - ET vs CE – more later on these two
- Can reduce somewhat with arrays of seismometers and subtraction of effect
- Forbiddingly large for $\sim 3\text{Hz}$ and lower
- **Ultimate limit on the lowest frequency detectors on- or under-ground**
- **And thus the largest BH masses detectable**



Thermal Noise

- **Thermal noise** – $k_B T$ of energy per mechanical mode
(*A. Einstein, 1905*)

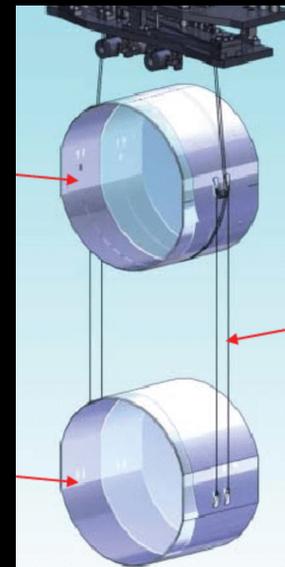
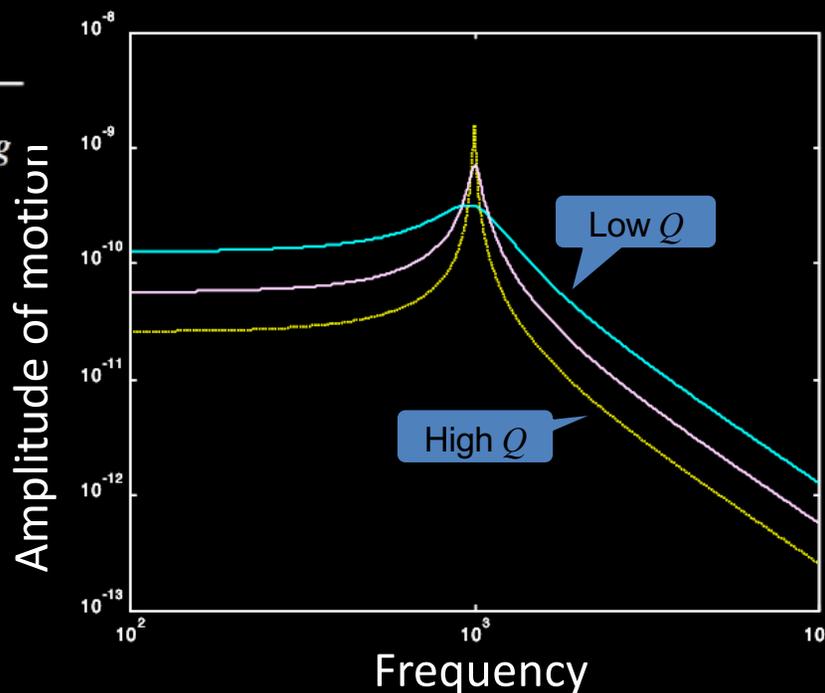
- Simple Harmonic Oscillator:

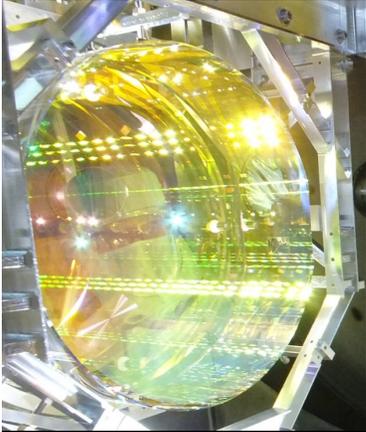
$$x_{rms} = \sqrt{\langle (\delta x)^2 \rangle} = \sqrt{k_B T / k_{spring}}$$

- Distributed in frequency according to real part of impedance $\Re(Z(f))$

$$\tilde{x}(f) = \frac{1}{\pi f} \sqrt{\frac{k_B T}{\Re(Z(f))}}$$

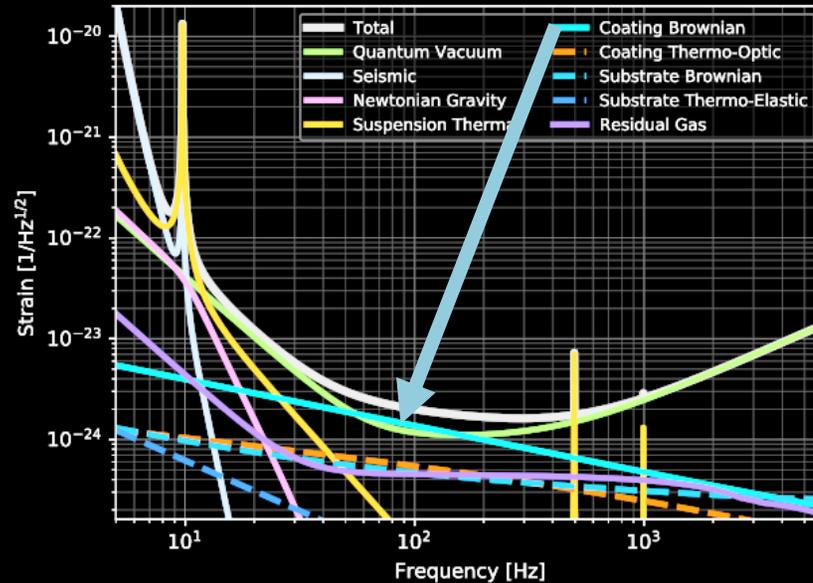
- **Low-loss materials**
- **monolithic construction**





Mirror Coating thermal noise

- In the best amorphous coatings, the dielectric optical coating has a rather large loss tangent
 - Some 10^{-4} , compared to 10^{-8} for fused silica
- The Fluctuation-Dissipation theorem says this is where the greatest motion is found
- And: the coating is the surface that is sensed by the laser
- **This is the dominant limit in the critical 50-200 Hz band**



coating thickness $\phi \equiv \text{Im}Y / \text{Re}Y$

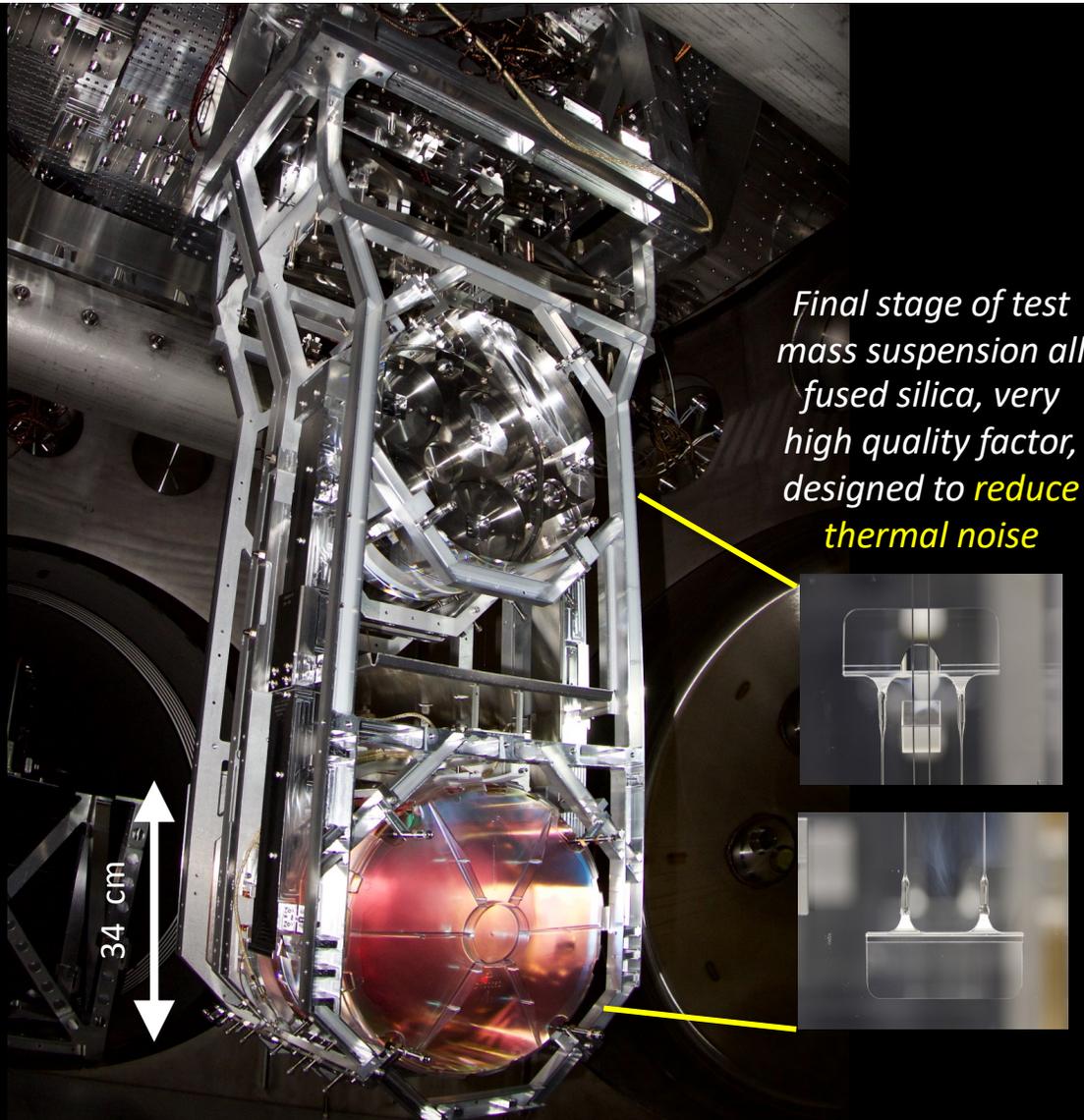
coating elastic loss

$$\langle \Delta x(f, T)^2 \rangle \approx \frac{2k_B T}{\pi^2 f} \frac{d}{w^2 Y} \phi(f)$$

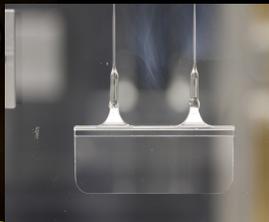
beam radius

Y Levin *Phys. Rev. D* **57** 659 (1998)

Test Mass Suspension

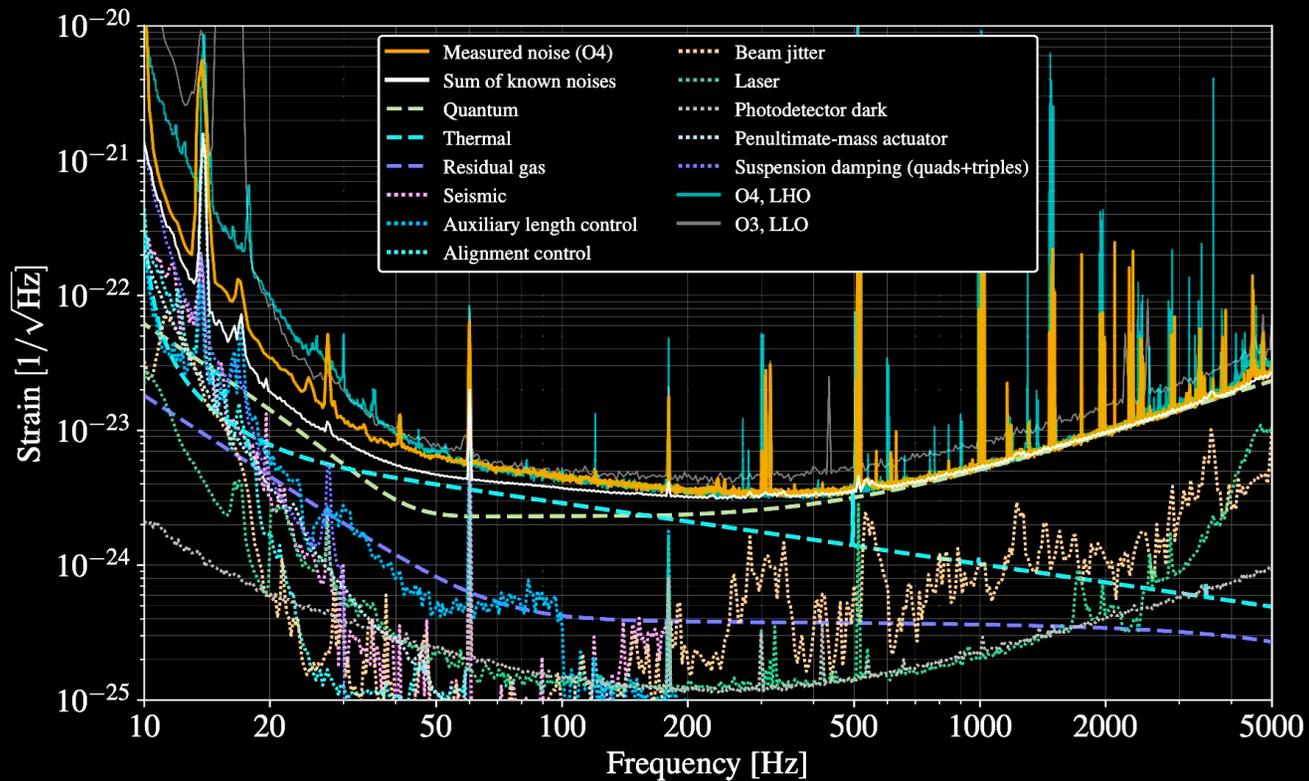


Final stage of test mass suspension all fused silica, very high quality factor, designed to **reduce thermal noise**



Rai Weiss,
~2020

LIGO 'Noise budget'



(b) Noise budget for the LIGO Livingston Observatory, as of October 2023.

Even better detectors would deliver more science. How to build a such a 10x detector?

- Current concept
- Limitations to the sensitivity
- **Scaling laws**

Sensing noises scale with arm length at various powers of $1/L$ – all get better

Shot Noise
while maintaining bandwidth

$$\frac{h_{\text{shot}}}{h_{0\text{shot}}} = \sqrt{\frac{2 \text{ MW}}{P_{\text{arm}}}} \sqrt{\frac{\lambda}{1.5 \mu\text{m}}} \left(\frac{3}{r_{\text{sqz}}} \right) \sqrt{\frac{40 \text{ km}}{L_{\text{arm}}}}$$

Radiation Pressure Noise
while maintaining bandwidth

$$\frac{h_{\text{RPN}}}{h_{0\text{RPN}}} = \sqrt{\frac{P_{\text{arm}}}{2 \text{ MW}}} \sqrt{\frac{1.5 \mu\text{m}}{\lambda}} \left(\frac{3}{r_{\text{sqz}}} \right) \left(\frac{40 \text{ km}}{L_{\text{arm}}} \right)^3$$

Coating Thermal Noise
loss angle dependence

$$\frac{h_{\text{CTN}}}{h_{0\text{CTN}}} = \sqrt{\frac{T}{123 \text{ K}}} \sqrt{\frac{\phi_{\text{eff}}(T)}{5 \times 10^{-5}}} \left(\frac{40 \text{ km}}{L_{\text{arm}}} \right)^{3/2}$$

Residual Gas Noise
facility limit

$$\frac{h_{\text{gas}}}{h_{0\text{gas}}} = \sqrt{\frac{p_{\text{gas}}}{4 \times 10^{-7} \text{ Pa}}} \sqrt{\frac{40 \text{ km}}{L_{\text{arm}}^{3/2}}}$$

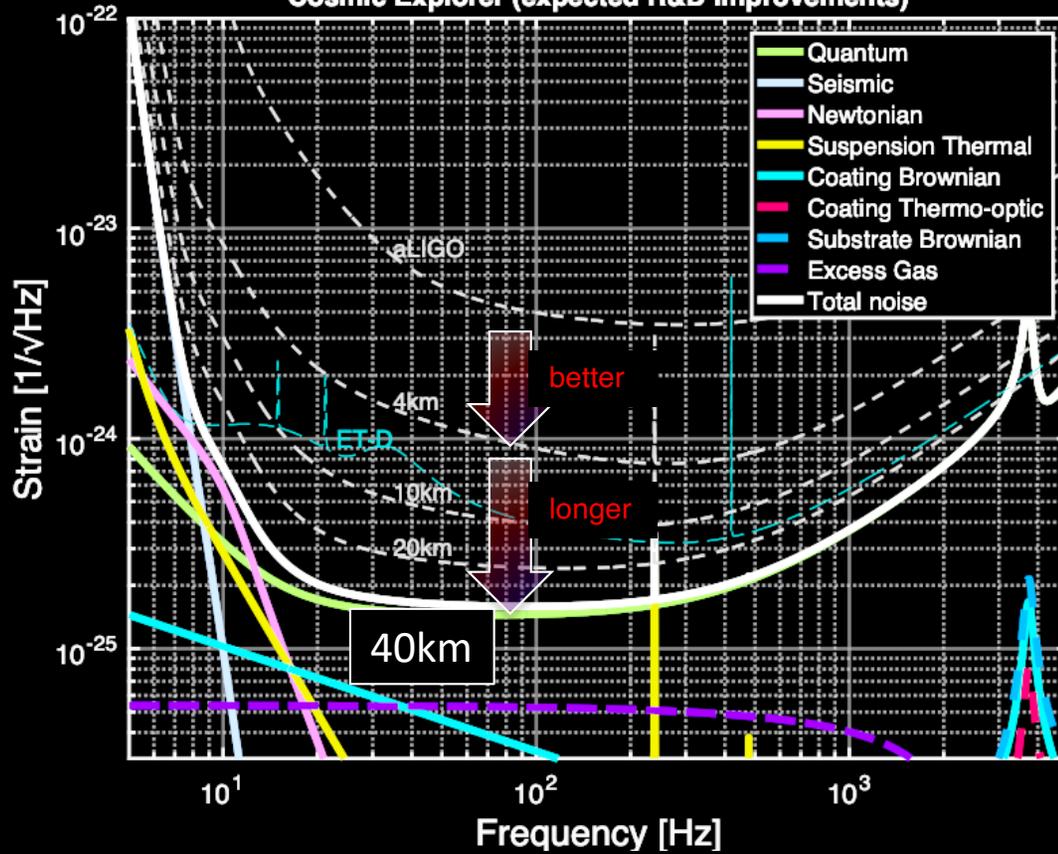
$$\Delta L = h L$$

Noise due to stochastic forces is independent of armlength

- Seismic noise; acoustic coupling
 - Varying direction of Newtonian gravity
 - Thermal noise motion (pendulum, substrate, coating)
- ...and the change in optical pathlength for a given GW-induced strain becomes larger, so SNR improves linearly with length
 - Up to $1/2 \lambda_{\text{GW}}$ – giving an optimal length for a given signal

Scaling Up... an example

Cosmic Explorer (expected R&D improvements)



Exploring the sensitivity of next generation
gravitational wave detectors
(2017) CQG 34, 044001

COSMIC EXPLORER

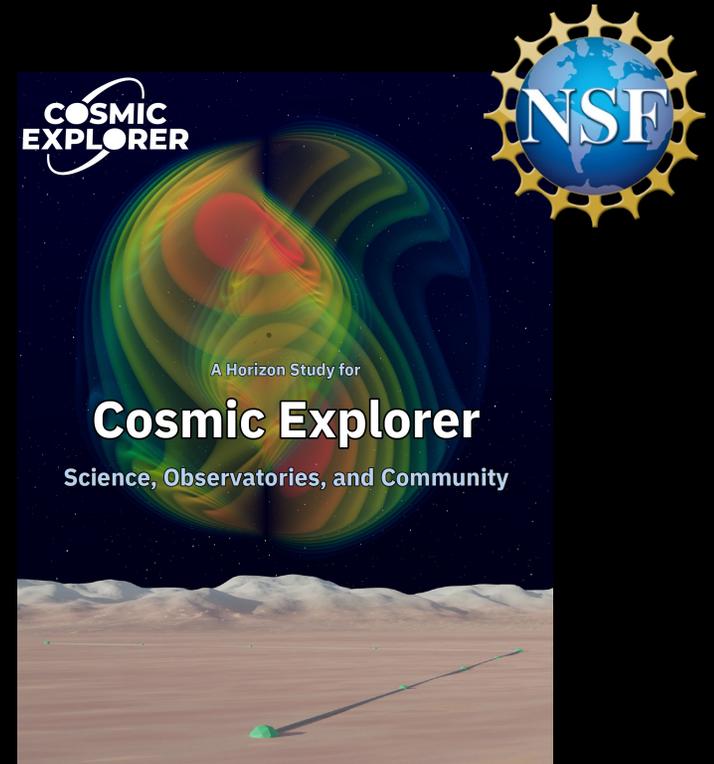
Astro 2020

Decadal Survey on Astronomy and Astrophysics

The National Academies of SCIENCES ENGINEERING MEDICINE



Image: Eddie Anaya (Cal State Fullerton)



- US contribution to the next-gen Network
- LIGO-like concept for a single interferometer per site, on Earth's surface
- CE is a larger, and more technically advanced version of LIGO: baseline of two widely separated observatories, 40km and 20km arms

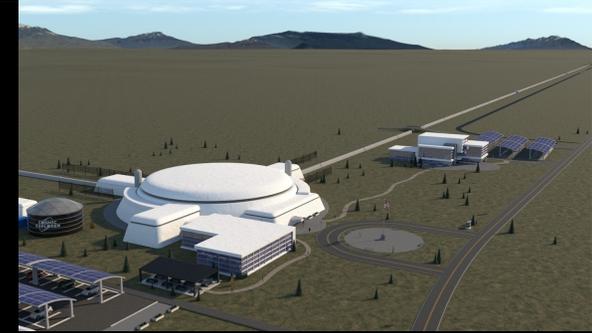
cosmicexplorer.org

CE Detector Design

- LIGO is starting to plan upgrades to the LIGO 4km detectors for the 'Post-O5' epoch – ~2029 to the start of CE observing (*optimistically* 2035)
- **Initial CE detectors will use all the techniques of this LIGO Post-O5 upgrade**
 - Move from a 4km baseline to a 40km baseline
 - **Low risk** – no significant advances in the detector are needed
 - Some work on bigger masses, suspensions, lower-loss optics
- Much like LIGO, later upgrades to CE are expected

CE Infrastructure

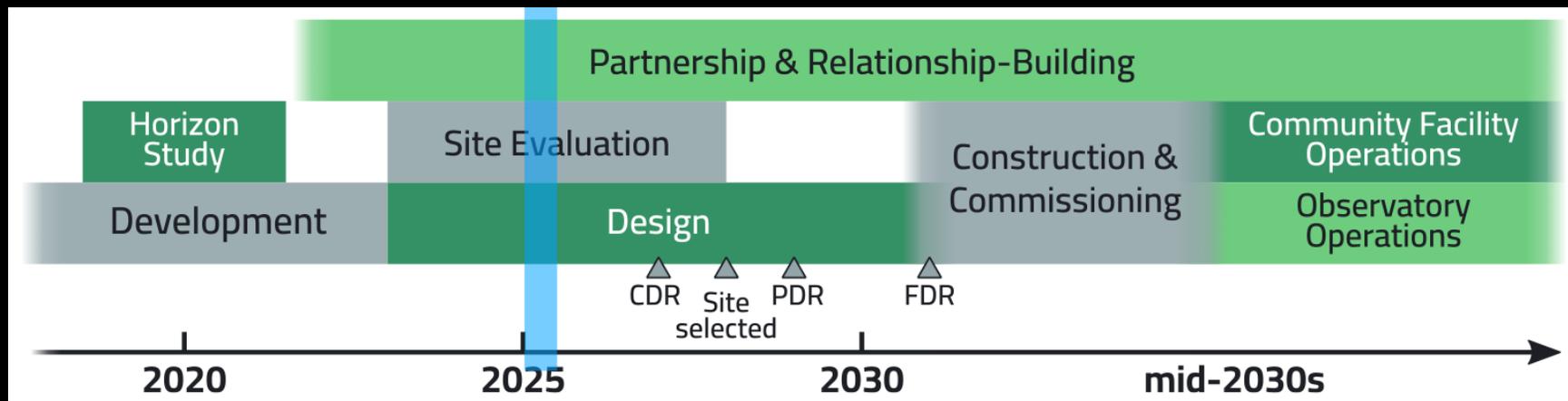
- Baseline of 40km and 20km observatories
 - 20km is ideal for observing $\sim 2\text{kHz}$ endgame of neutron-star mergers
 - 40km is optimized for absolute 'reach' to enable all in-scope binaries to be observed
- Sites separated by a continental baseline
 - Hope that ET in Europe will be built; expect that data will be shared
- Working on less expensive vacuum systems (dominates the cost)
- Single interferometric detector per site
- Earth's surface construction
 - Good: less expensive than tunneling; no complexity of underground work; future modifications of interferometer layout easier (no new caverns)
 - Bad: increased coupling to surface 'seismic' noise, and thus Newtonian background – limits low-frequency sensitivity ($\sim 7\text{Hz}$ compared to $\sim 5\text{Hz}$ for ET)
- *Geographically* suitable sites can be found in the US (and Canada, Australia...)



CE Status

- Conceptual Design is now underway
- CE Funding is being sought for this phase
 - Working with the NSF; their model is multiple smaller awards
 - International contributions (in-kind) – UK, Germany currently
- CE Project Phase Funding
 - Order of $\$2 \times 10^9$ for the two-detector baseline
 - Hope that the NSF continues to see GWs as a compelling topic
 - DOE, private funding, in-kind or other international support all sought
- Key participant in the future network: CE, ET, LIGO-India

Cosmic Explorer Timeline ...assuming funding



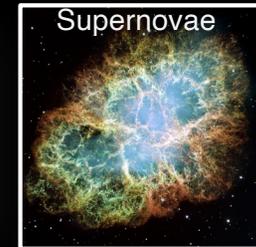
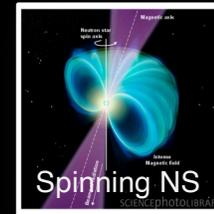
We are here

The last page (at last!)

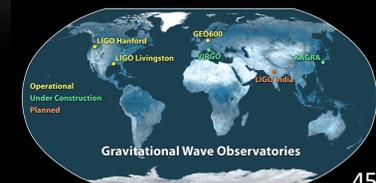
- Ground-based GW observation works
- There are lots of sources yet to be observed
- Scaling laws show the technical feasibility of better detectors
- The US Concept, Cosmic Explorer is our proposal for a Next-Gen GW detector
 - Top-down observational science prioritization
 - bottom-up technology study
- The biggest challenges appear to be in
 - finding resources to build Cosmic Explorer
 - building relationships with host communities
- **CE needs very broad community support to merit the investment**

Why stop there?

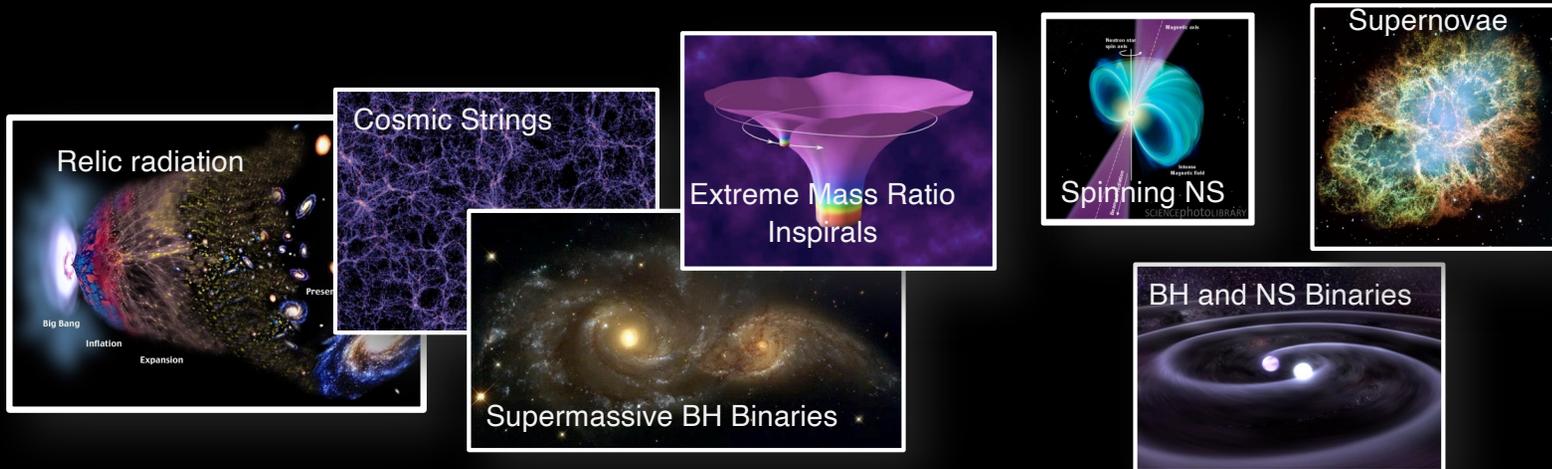
The Gravitational Wave Spectrum



Terrestrial interferometers



The Gravitational Wave Spectrum



10^{-16} Hz

10^{-9} Hz

10^{-4} Hz

10^0 Hz

10^3 Hz

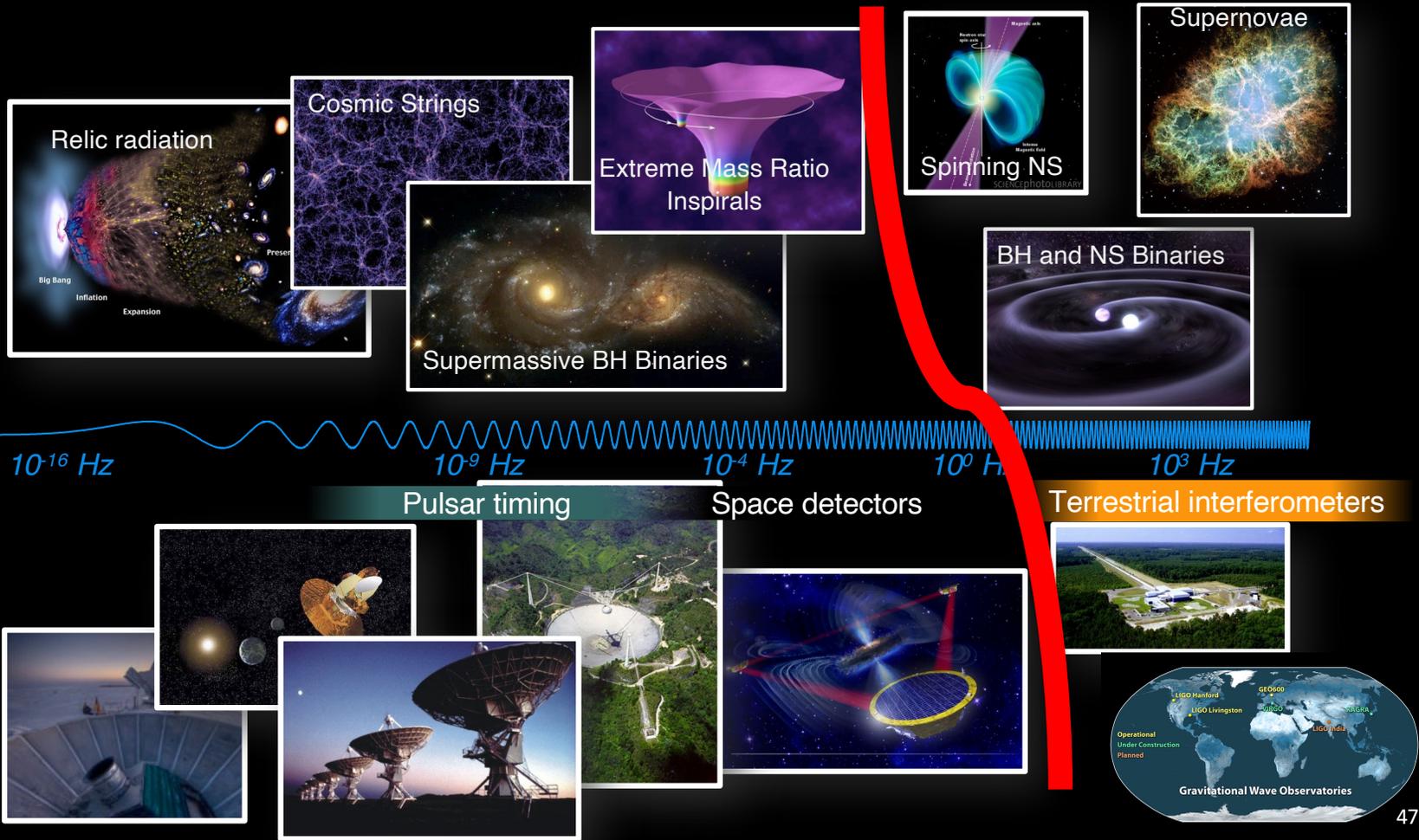
Pulsar timing

Space detectors

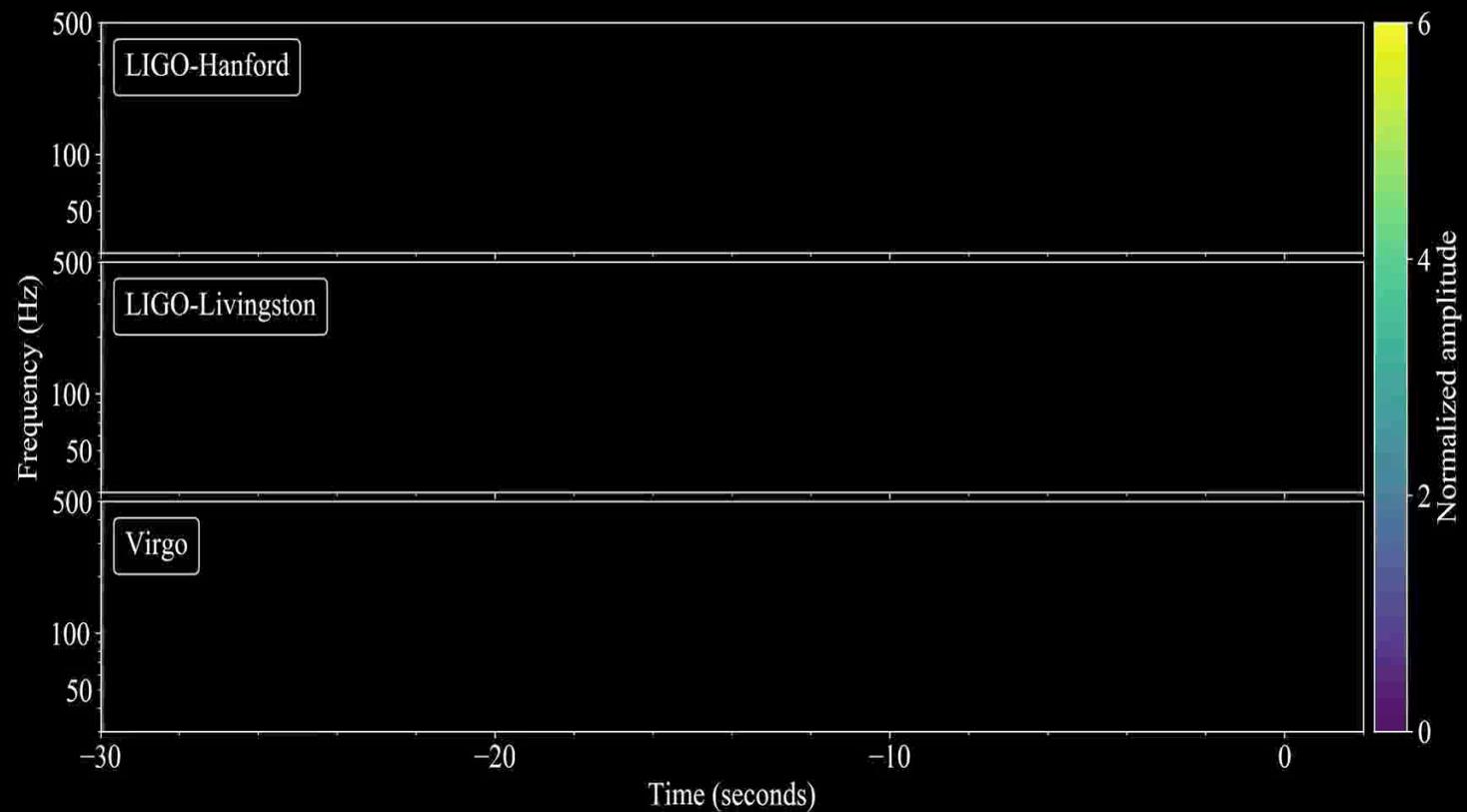
Terrestrial interferometers



The Gravitational Wave Spectrum



The binary neutron star signal, with and without the interferometer noise



Einstein Telescope

Underground; in Sardinia or near NL-B-D border

Six detectors arranged in a **triangle**

⇒ sense both GW polarizations

⇒ “null stream” for consistency checks

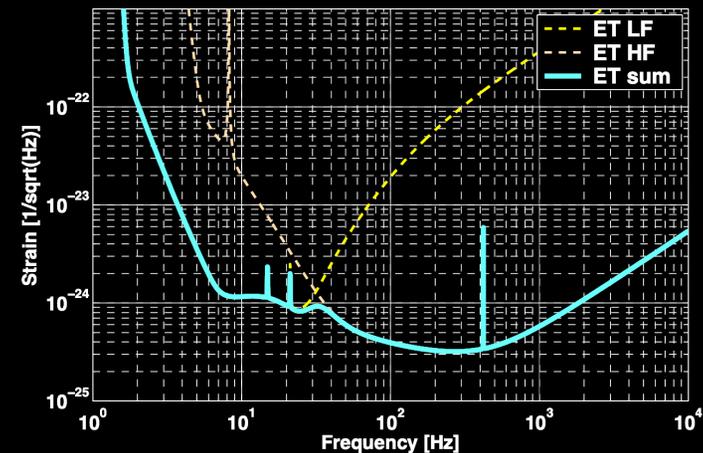
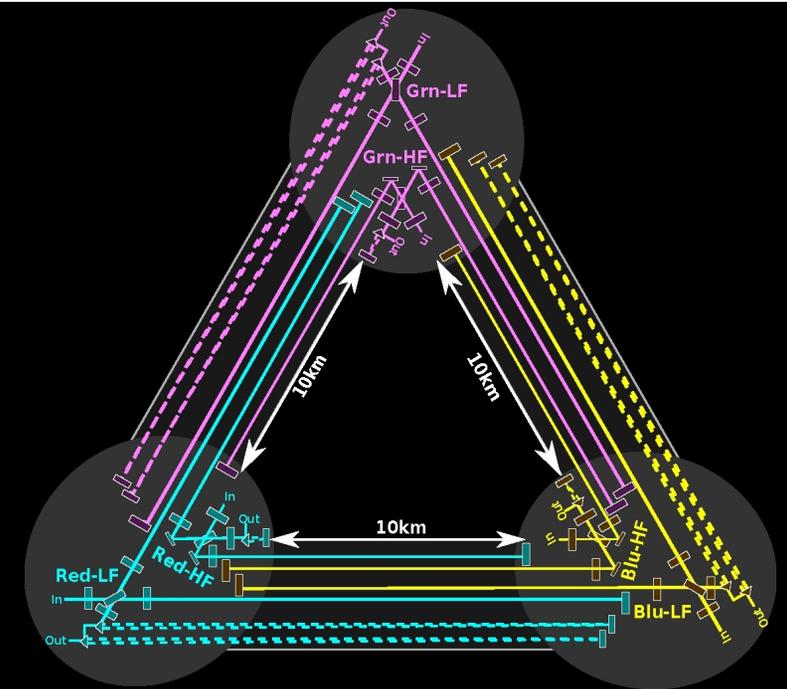
Observe down to a few hertz and have deep bucket using “**xylophone**” detectors with large masses, freq-dep squeezing

⇒ low frequency: cryogenic silicon

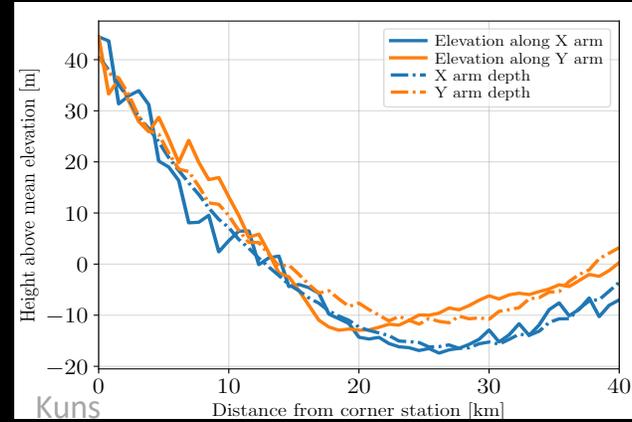
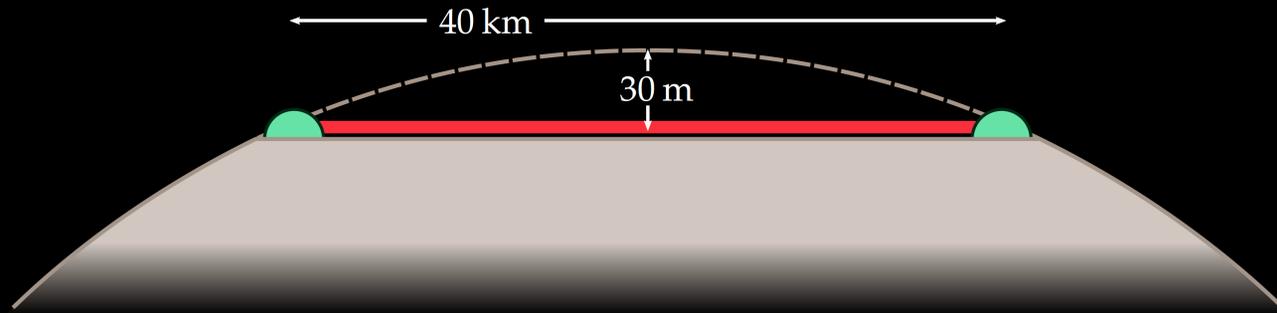
⇒ high frequency: high power, room temp

ET Design Report published 2011, [updated](#) in 2020

In 2021, ET was [included in the roadmap](#) of the European Strategic Forum for Research Infrastructures (ESFRI).



40km CE



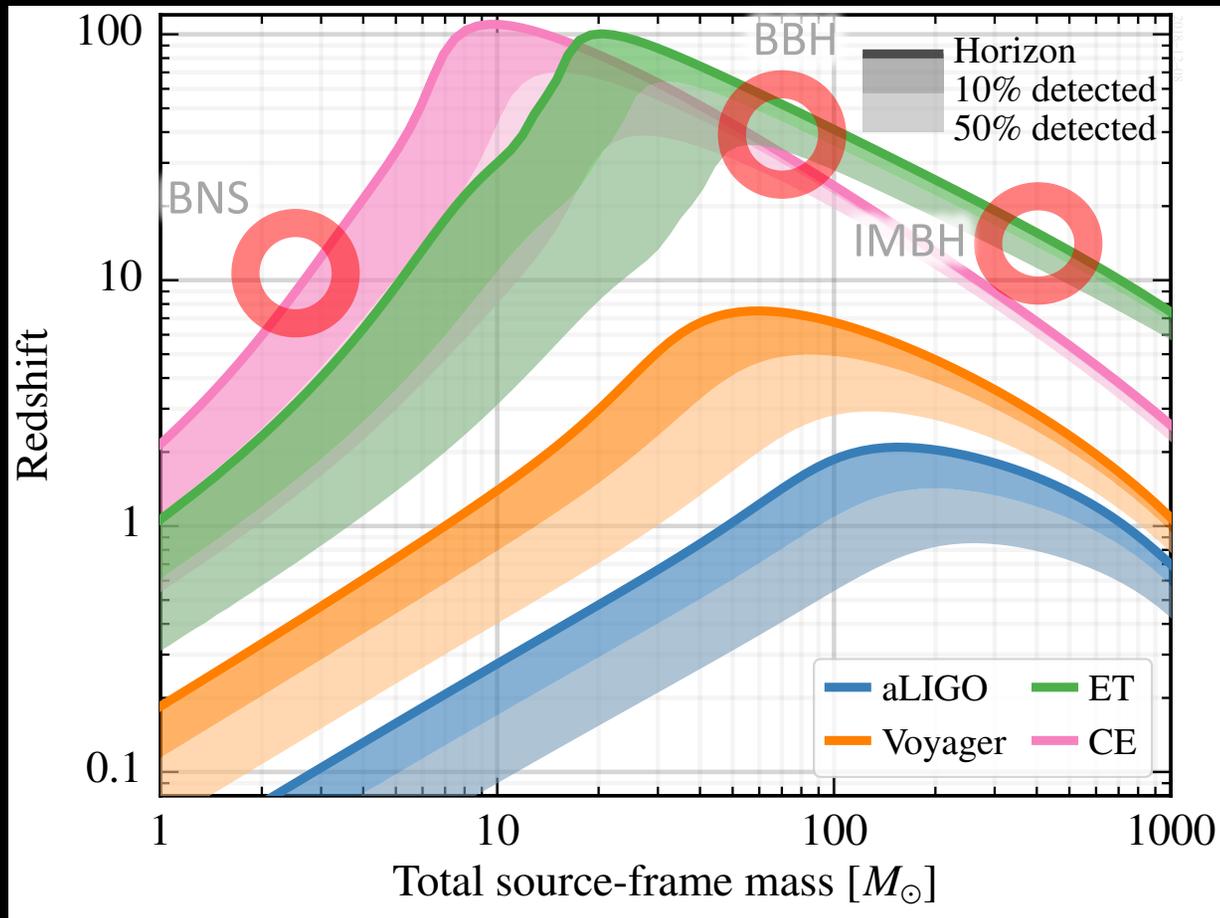
Why 40km?

- Broadly speaking, the sensitivity of these instruments improves with length
- The bandwidth is, however, limited to roughly

$$\frac{c}{2L} = \frac{3 \times 10^5 \frac{km}{s}}{2 \times 40 km} \simeq 4 kHz$$

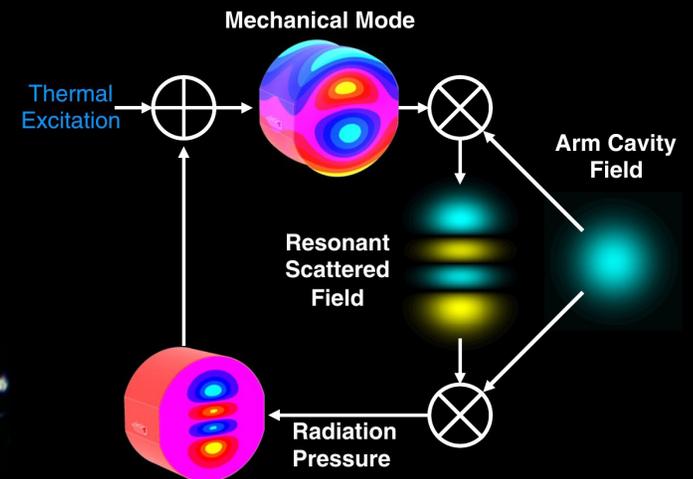
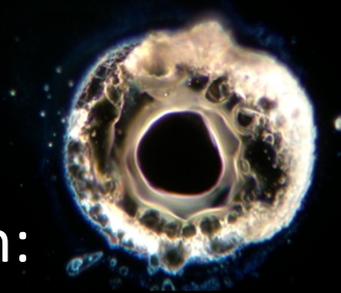
so making a detector longer than 40km would compromise its access to interesting astrophysics (i.e., post-merger signals and supernovae).

What can CE do?



Noise improvements: reducing quantum noise

- Increasing the laser power in the arms
O1,O2 (100kW) → O3 (200kW) → goal is 400 kW for O4
- Not easy!
 - You need a high power laser first..
 - Mirror radii must remain within a few meters of the ~2 kilometer nominal value
 - Control issues: angular control and parametric instabilities
 - “Point absorbers”
[Applied Optics Vol. 60, Issue 13 pp. 4047-4063 \(2021\)](#)



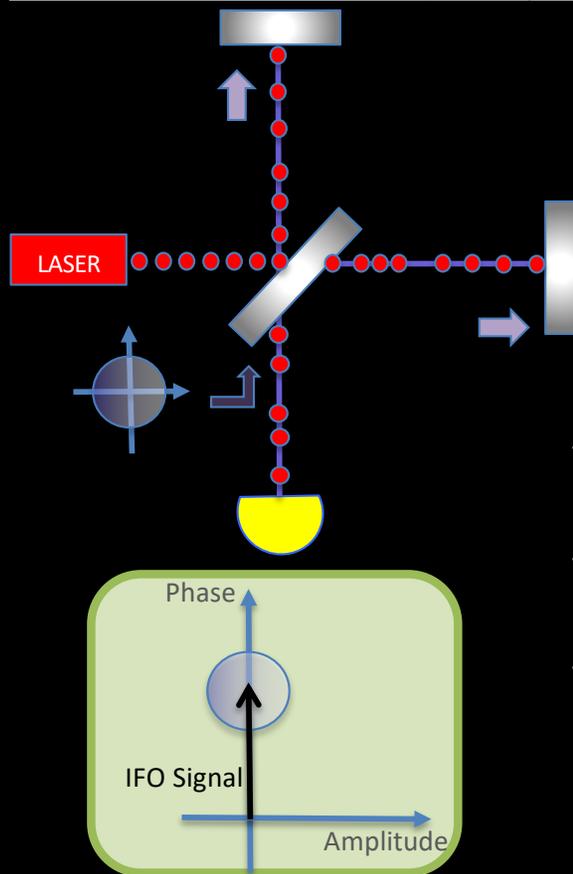
- Complementary approach:
squeezed states of vacuum

Quantum-mechanical noise in an interferometer

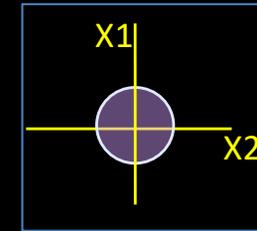
Carlton M. Caves

W. K. Kellogg Radiation Laboratory, California Institute of Technology, Pasadena, California 91125

(Received 15 August 1980)



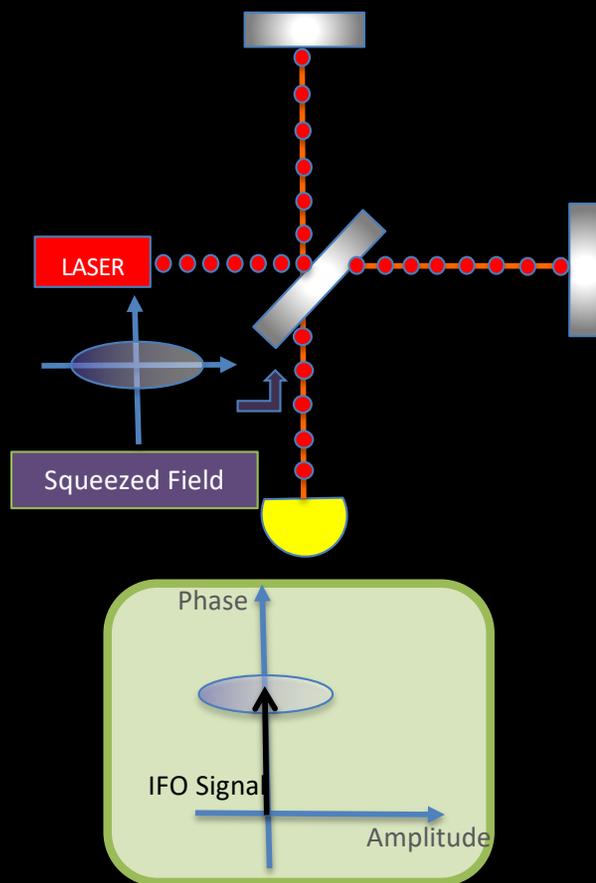
Zero-point energy
(vacuum) fluctuations



$$\Delta X_1 \Delta X_2 \geq 1$$

- ✧ When average amplitude is zero, the variance remains
- ✧ Heisenberg uncertainty principle, quadratures associated with **amplitude** and **phase**
- ✧ They enter the interferometer from all the open ports of the interferometer, but the ones which matter are the one **entering from the anti-symmetric port!**

Replace regular vacuum with squeezed vacuum

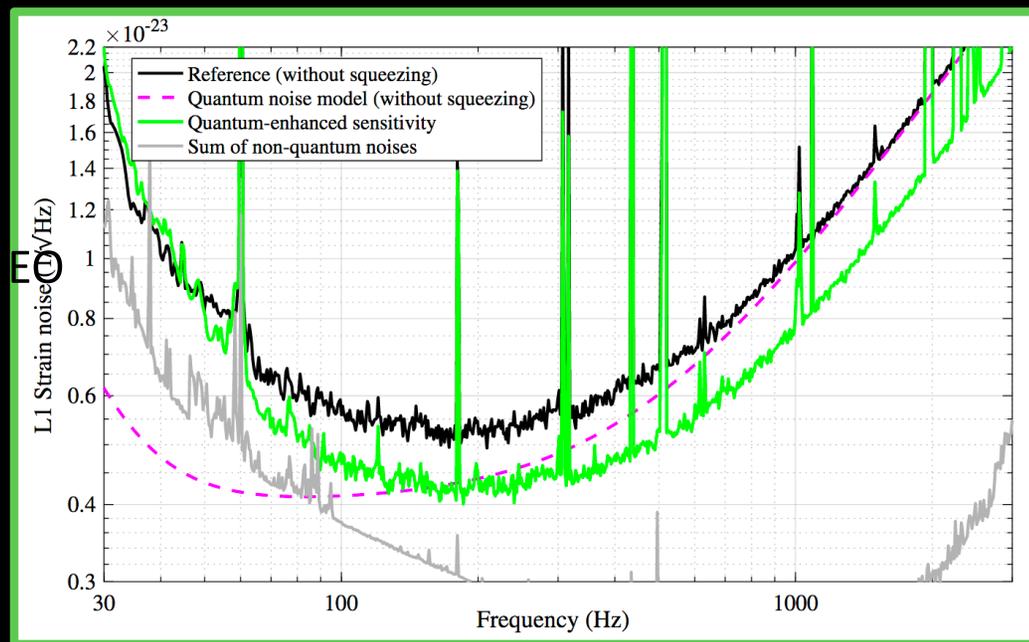


- ✧ Reduce quantum noise by injecting squeezed vacuum: less uncertainty in one of the two quadratures
- ✧ Heisenberg uncertainty principle: if the noise gets smaller in one quadrature, it gets bigger in the other one
- ✧ One can choose the relative orientation between the squeezed vacuum and the interferometer signal (squeeze angle)
- ✧ Squeezing is made by creating pairs of photons using an optical parametric oscillator
- ✧ The pairs are quantum-mechanically entangled and have correlated arrival times at the detector
- ✧ This reduces the randomness of the time distribution

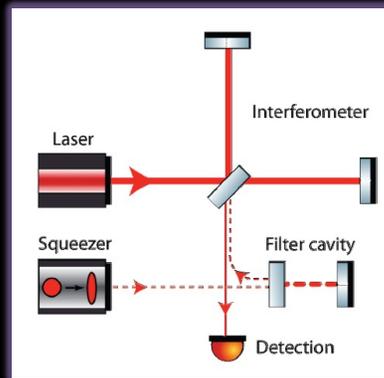
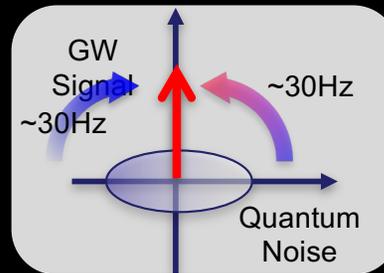
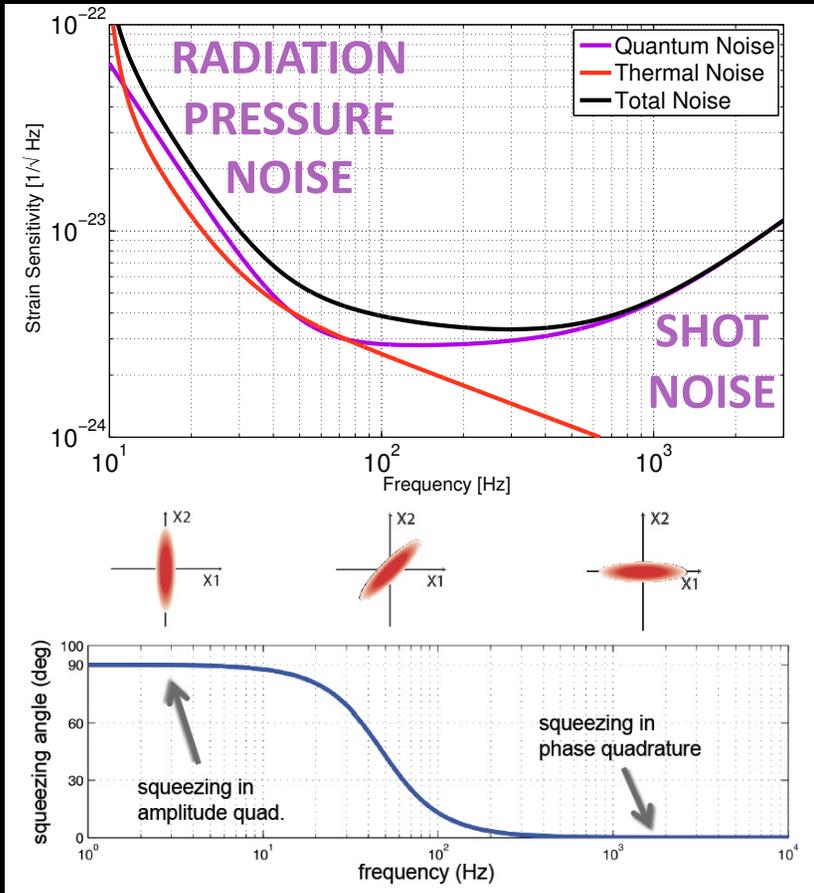
Squeezing performance in O3

[PhysRevLett.123.231107](#) *Nature* 583, pages 43–47 (2020)

3 dB of squeezing observed at high frequency = 40% quantum noise reduction (in amplitude); observation of quantum radiation pressure noise in both detectors



Frequency Dependent Squeezing for O4



High finesse detuned “filter cavity” which rotates the squeezing angle as function of frequency



Credit: Antonio Pasqualetti

Highlight from Virgo:
300 m filter cavity already built and locked and characterized, commissioning in progress 58